

Relativistic Evaluation of the Charge-Transfer Probabilities and Cross Sections for Low-energy Collisions of H-like Ions with Bare Nuclei

[Yury Kozhedub](#)

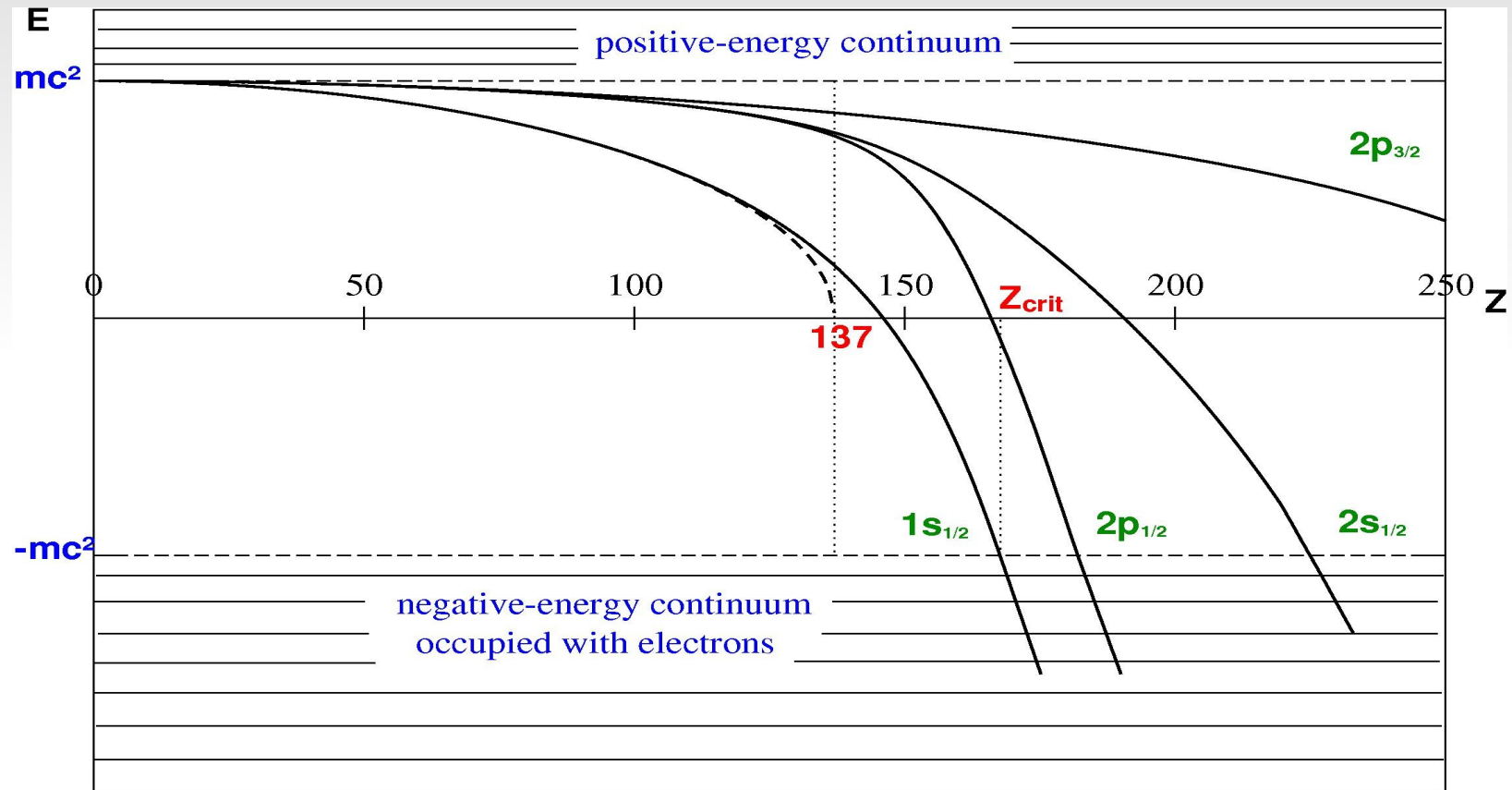
St. Petersburg State University

Outline:

- Introduction and motivation
- Theoretical description
- Numerical results
- Summary

Introduction

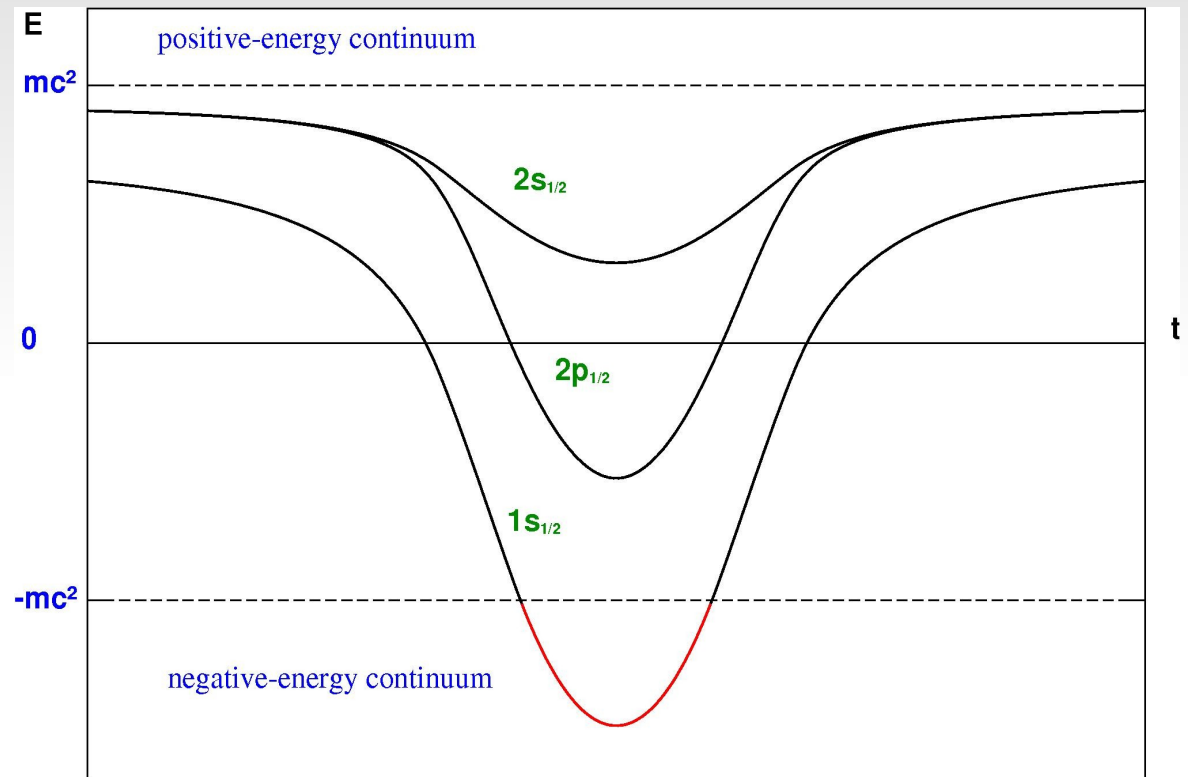
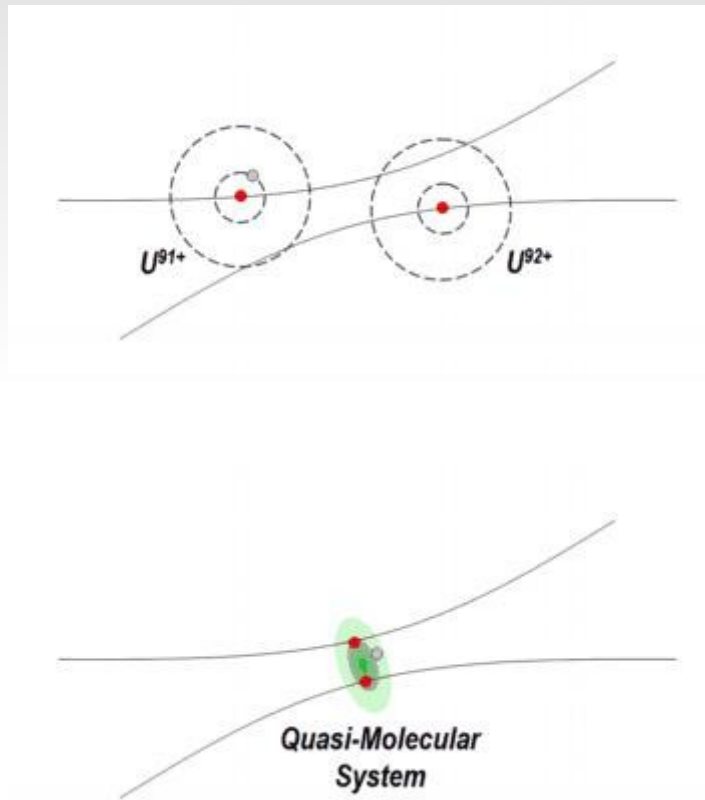
The $1s$ level dives into the negative-energy continuum at $Z_{\text{crit}} \sim 173$.



[S.S. Gershtein, Ya.B. Zeldovich, 1969; W. Pieper, W. Greiner, 1969]

Motivation: super-heavy quasi-molecules

A way to create super-critical field is to collide two heavy ions with $Z_1 + Z_2 > 173$.



Two center Dirac equation

Features of the investigated process:

- Low-energy ions: $\sim 6 \text{ MeV/u}$ for U^{91+} - U^{92+}
- Relativistic electron: $v_e \sim (\alpha Z)c$
- $m_e \ll M_{\text{nucl}} \rightarrow$ Nuclei ($\mathbf{R}_A, \mathbf{R}_B$) move according to the Rutherford trajectory

The time-dependent and stationary (for fixed R_{AB}) Dirac equations (in a.u.)

$$i \frac{d\Psi}{dt} = h_D \Psi(\vec{r}, t), \quad h_D \psi_n(\vec{r}) = \varepsilon_n \psi_n(\vec{r}),$$

$$h_D = c(\vec{\alpha} \cdot \vec{p}) + \beta mc^2 + V_{AB}(\vec{r}),$$

where $\vec{\alpha}, \beta$ are the Dirac matrices, and $V_{AB}(\vec{r}) = V_{\text{nucl}}^{(A)}(\vec{r}_A) + V_{\text{nucl}}^{(B)}(\vec{r}_B)$,

$$\vec{r}_A = \vec{r} - \vec{R}_A, \quad \vec{r}_B = \vec{r} - \vec{R}_B.$$

Finite Basis Expansion

$$\Psi(\vec{r}) = \sum_i c_i \varphi_i(\vec{r}),$$

$$\Psi(\vec{r}, t) = \sum_i C_i(t) \varphi_i(\vec{r}).$$

Stationary case:
$$\sum_j S_{ij} c_j = \sum_j H_{ij} c_j$$

Time-dependent case:
$$i \sum_j S_{ij} \frac{dC_j(t)}{dt} = \sum_j (H_{ij} - T_{ij}) C_j(t),$$

where

$$H_{ij} = \langle \varphi_i | h_D | \varphi_j \rangle, \quad T_{ij} = i \langle \varphi_i | \frac{\partial}{\partial t} | \varphi_j \rangle, \quad S_{ij} = \langle \varphi_i | \varphi_j \rangle.$$

Basis set

- Our basis is constructed as a sum of the Dirac and Dirac-Sturm orbitals, localized on each ion.
- The Dirac and Dirac-Sturm orbitals are obtained by solving numerically the finite-difference one-center Dirac and Dirac-Sturm equations.

$$\Psi(\vec{r}, t) = \sum_{\alpha=A, B} \sum_{\mu} C_{\alpha, \mu}(t) \varphi_{\alpha, \mu}(\vec{r} - \vec{R}_{\alpha}(t))$$

$\varphi_{\alpha, \mu}$ is the Dirac or the Dirac-Sturm orbital localized on the center α .

Basis set advantages

- Spectrum of the Dirac-Sturm operator is red **discrete** and **complete** (including functions of the **negative** Dirac spectrum).
- Relativistic DSO satisfy **the dual kinetic balance condition** [*V. Shabaev et al., PRL 93, 130405 (2004)*].
- DSO have correct **asymptotic behavior** when $r \rightarrow 0$ and $r \rightarrow \infty$.
- All DSO have approximately **the same space scale**, which does not depend on the principal quantum number n .
- **Monopole approximation** enables partly accounting for the potential of the second ion in constructing of the basis functions.

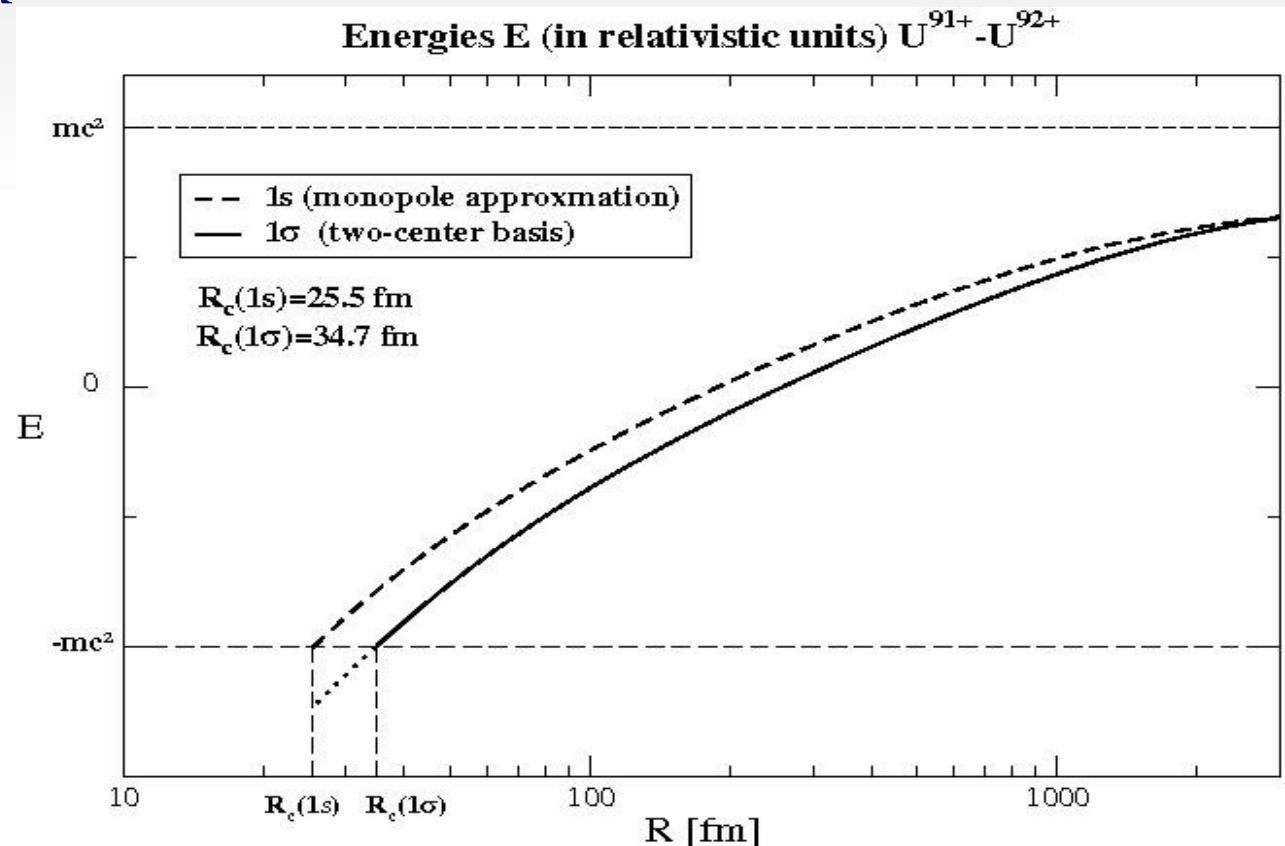
The Basis set

- Provides the natural satisfaction of **the initial conditions**.
- Allows one to evaluate **the ionization cross section**.
- Is perfect for describing the quasi-molecular states at **small inter-nuclear distance**. This is especially important for investigation of the diving effect.
- Possesses **fast basis convergence**, that significantly reduces the size of matrix problem and calculation time.

Energies of the $1\sigma_+$ ground states of quasi-molecules

The $1\sigma_+$ state energy of the U_2^{183+} quasi-molecule as a function of the internuclear distance R

$$R(1\sigma_+) = 34.7 \text{ fm}$$



Critical Distances R_c (fm)

Z	Point nucleus		Extended nucleus	
	This work	Others	This work	Others
88	24.27	24.24 ^a	19.91	19.4 ^d
90	30.96	30.96 ^a	27.06	26.5 ^d
92	38.43	38.4 ^b	34.74	34.7 ^b
		38.42 ^a		34.3 ^d
		36.8 ^c		34.7 ^f
94	46.58	46.57 ^a	43.13	42.6 ^d
96	55.38	55.37 ^a	52.10	
98	64.79	64.79 ^a	61.61	61.0 ^d
				61.1 ^f

^a [V.I. Lisin et al., *Phys. Lett.* **69B**, 2 (1977)]

^b [A. Artemyev et al., to be published]

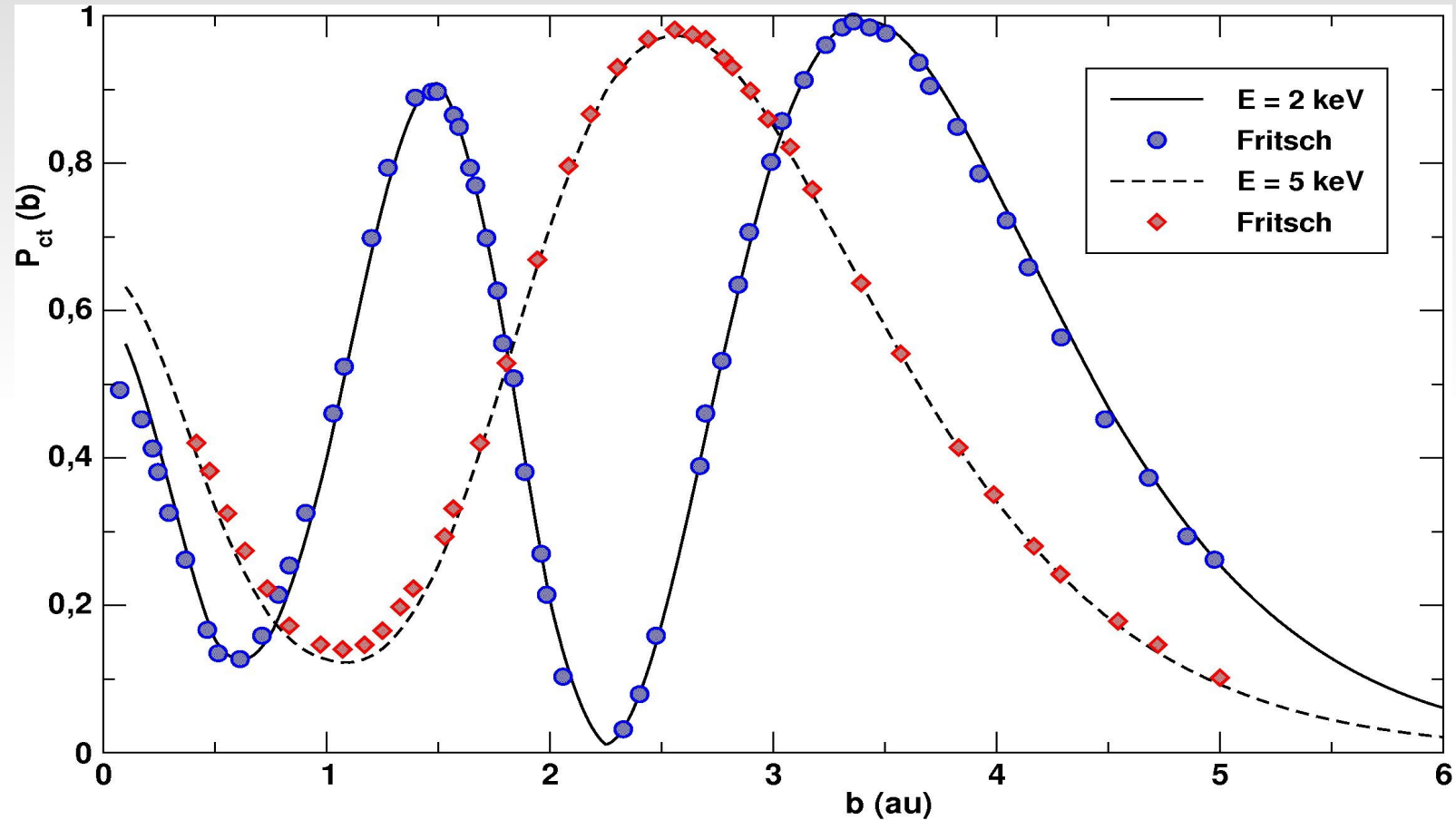
^c [J. Rafelski, B. Müller, *Phys. Lett.* **65B**, 205 (1976)]

^d [V.I. Lisin et al., *Phys. Lett.* **91B**, 20 (1980)]

^f [B. Müller and W. Greiner, *Z. Naturforsch.* **31a**, 1 (1975)]

H(1s)-H⁺ collision

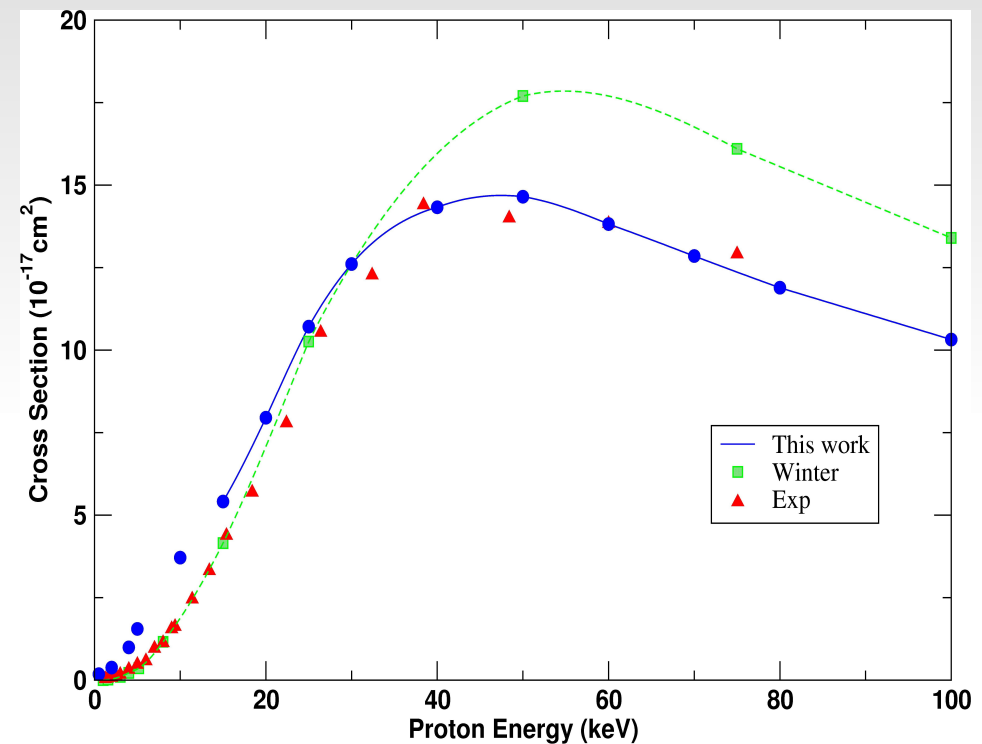
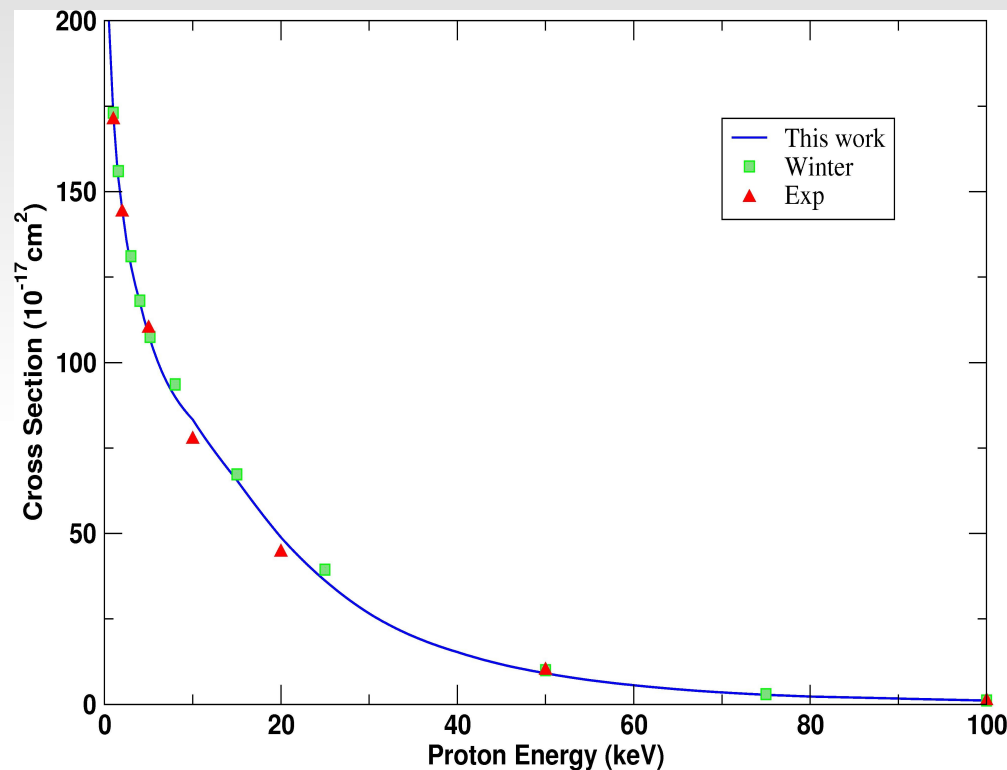
Charge-transfer probability as a function of the impact parameter b



Comparison with the results of work [*W. Fritsch et al., Phys. Rep., 202, 1 (1991)*]

H(1s)-H⁺ collision

Charge-transfer and Ionization cross sections as functions of the collision energy



Other calculations [G. Winter, *PRA* **80**, 032701 (2009)]

Experimental data [R. Janev et al., *At. and Pl. Mat. Int. Data for Fusion*, Nucl. Fusion Suppl 4 (1993); M. Shah et al., *JPB* **31**, L757 (1998); **20**, 2481 (1987)]

Z Scaling Transformation

In the nonrelativistic case, using the scale transformation it is possible to establish a link between the charge-transfer parameters in the $H(1s)-H^+$ and $A^{(Z-1)+}(1s)-A^{Z+}$ collisions

$$\begin{aligned}\vec{r}' &= Z \vec{r} \\ \vec{R}'_A &= Z \vec{R}_A \\ \vec{R}'_B &= Z \vec{R}_B \\ t' &= Z^2 t \\ E' &= E / Z^2 \\ \vec{V}'_A &= \vec{V}_A / Z \\ \vec{V}'_B &= \vec{V}_B / Z \\ \sigma' &= \sigma Z\end{aligned}$$

Ne¹⁰⁺(1s)-Ne⁹⁺ collision

Charge-transfer cross section σ_{ct} (10^{-17} cm²) for the Ne¹⁰⁺(1s)-Ne⁹⁺ and H(1s)-H⁺ collisions

Ne ¹⁰⁺ (1s)-Ne ⁹⁺				H(1s)-H ⁺
Energy <i>E/Z</i> ² (keV/u)	$\sigma_{\text{ct}} \cdot Z^2$ Rel.	$\sigma_{\text{ct}} \cdot Z^2$ Nonrel.	$\sigma_{\text{ct}} \cdot Z^2$ Born approx. ^a	$\sigma_{\text{ct}} \cdot Z^2$
1.0	171.6	172.2	188.4	172.4
4.0	117.1	117.5	114.8	117.5
10.0	80.8	81.3	76.2	81.3
20.0	48.5	48.9	48.2	48.9

^a *V. Shevelko, 2010*

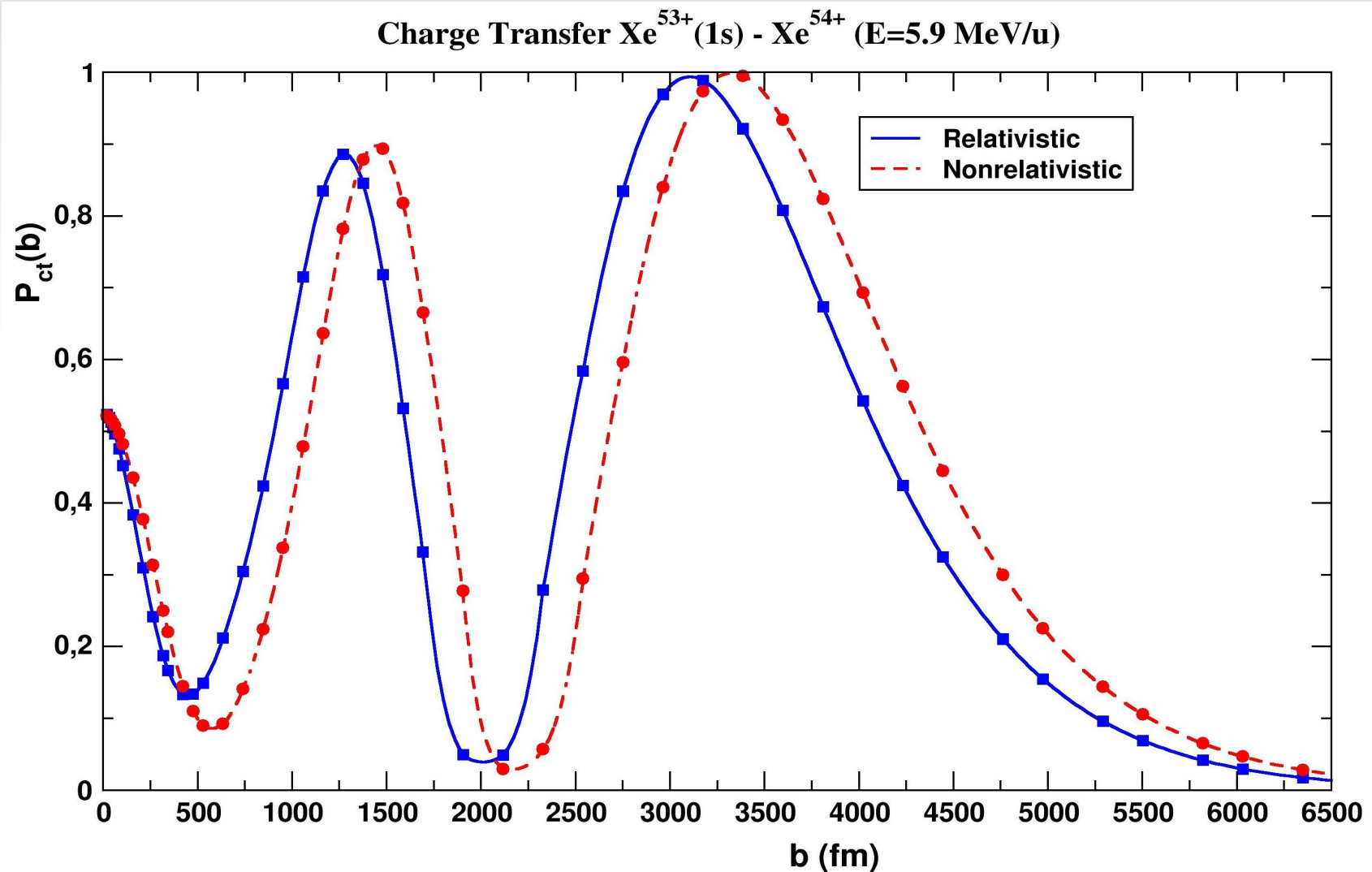
$\text{Xe}^{53+}(1s)\text{-Xe}^{54+}$ collision

Charge-transfer cross section σ_{ct} (10^{-17} cm^2) for the $\text{Xe}^{53+}(1s)\text{-Xe}^{54+}$ and $\text{H}(1s)\text{-H}^+$ collisions

	$\text{Xe}^{54+}(1s)\text{-Xe}^{53+}$			$\text{H}(1s)\text{-H}^+$
E/Z^2 (keV/u)	E (MeV/u)	$\sigma_{\text{ct}} \cdot Z^2$ Rel.	$\sigma_{\text{ct}} \cdot Z^2$ Nonrel.	$\sigma_{\text{ct}} \cdot Z^2$
1.234 57	3.6	148.3	163.3	165.0
2.023 32	5.9	129.4	143.0	144.9
3.429 36	10.0	109.1	123.8	124.8
34.293 6	100.0	13.3	20.6	20.7

$\text{Xe}^{53+}(1s)\text{-Xe}^{54+}$ collision

Charge-transfer probability as a function of the impact parameter b ,
 $E=5.9 \text{ MeV/u}$



Summary

Conclusion:

- A new method employing the Dirac-Sturm basis functions for evaluation of various electron-excitation processes in low-energy heavy-ion collisions has been developed
- Relativistic calculations of the charge transfer for low-energy collisions of H-like ions with bare nuclei have been carried out

Outlook:

- Calculations of the charge-transfer probabilities at supercritical regime, $Z_1 + Z_2 > 173$
- Investigation of a possibility of indirect observation of the diving effect

Thank You for Your Attention