

MATTER @ HIGH DENSITIES IN COMPACT STARS & CBM

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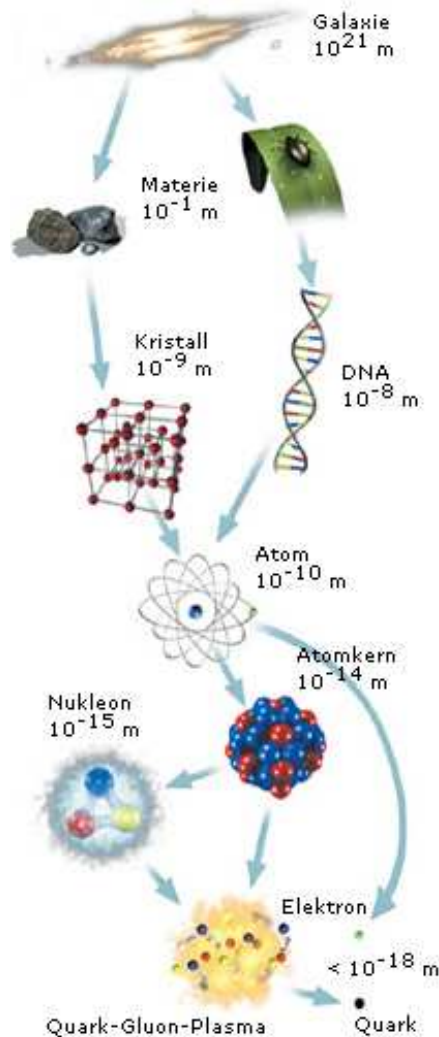
Institut für Physik, Universität Rostock, Germany

Bogoliubov Laboratory for Theoretical Physics, JINR Dubna, Russia

- **Introduction:** Many-particle Systems and Quantum Field Theory
- **Partition function for QCD:** Lattice Simulations vs. Resonance Gas
- **Bound states and Mott effect, Color superconductivity**
- **Application 1:** J/ψ suppression in Heavy-Ion Collisions
- **Application 2:** Quark Matter in Compact Stars and in CBM @ FAIR
 - Masses and Cooling behaviour
 - Flow constraint and phase diagram
- **Conclusions / Outlook**



MANY PARTICLE SYSTEMS & QUANTUM FIELD THEORY



Elements

Bound states

System

humans, animals

couples, groups, parties

society

molecules, crystals

(bio)polymers

animals, plants

atoms

molecules, clusters, crystals

solids, liquids, ...

ions, electrons

atoms

plasmas

nucleons, mesons

nuclei

nuclear matter

quarks, anti-quarks

nucleons, mesons

quark matter

Highly Compressed Matter \Leftrightarrow Pauli Principle

$$\text{Partition function: } Z = \text{Tr} \left\{ e^{-\beta(H - \mu_i Q_i)} \right\}$$

PARTITION FUNCTION FOR QUANTUM CHROMODYNAMICS (QCD)

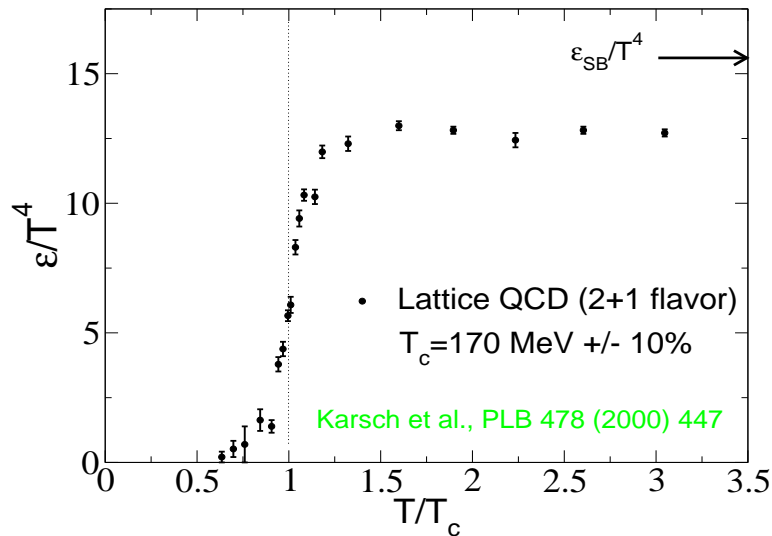
- Partition function as a Path Integral (imaginary time $\tau = i t, 0 \leq \tau \leq \beta = 1/T$) \Rightarrow PS I

$$Z[T, V, \mu] = \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \mathcal{D}A \exp \left\{ - \int_0^\beta d\tau \int_V d^3x \mathcal{L}_{QCD}(\psi, \bar{\psi}, A) \right\}$$

- QCD Lagrangian, non-Abelian gluon field strength: $F_{\mu\nu}^a(A) = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g f^{abc} [A_\mu^b, A_\nu^c]$

$$\mathcal{L}_{QCD}(\psi, \bar{\psi}, A) = \bar{\psi} [i\gamma^\mu (\partial_\mu - igA_\mu) - m - \gamma^0 \mu] \psi - \frac{1}{4} F_{\mu\nu}^a(A) F^{a,\mu\nu}(A)$$

- Numerical evaluation: Lattice gauge theory simulations (Bielefeld group)



- Equation of state: $\epsilon(T) = -\partial \ln Z[T, V, \mu] / \partial \beta$

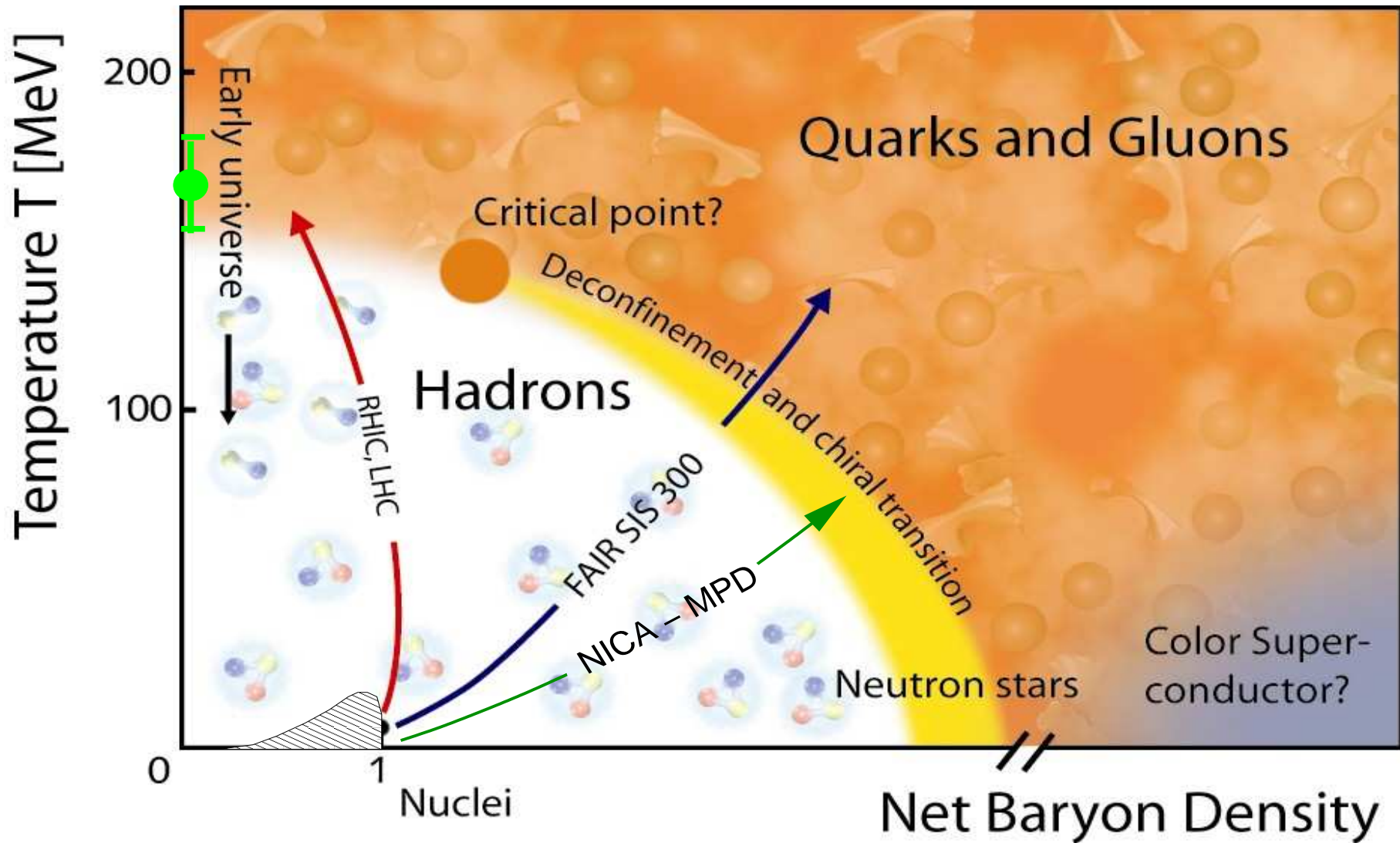
- Phase transition at $T_c = 170$ MeV

- **Problem:** Interpretation ?

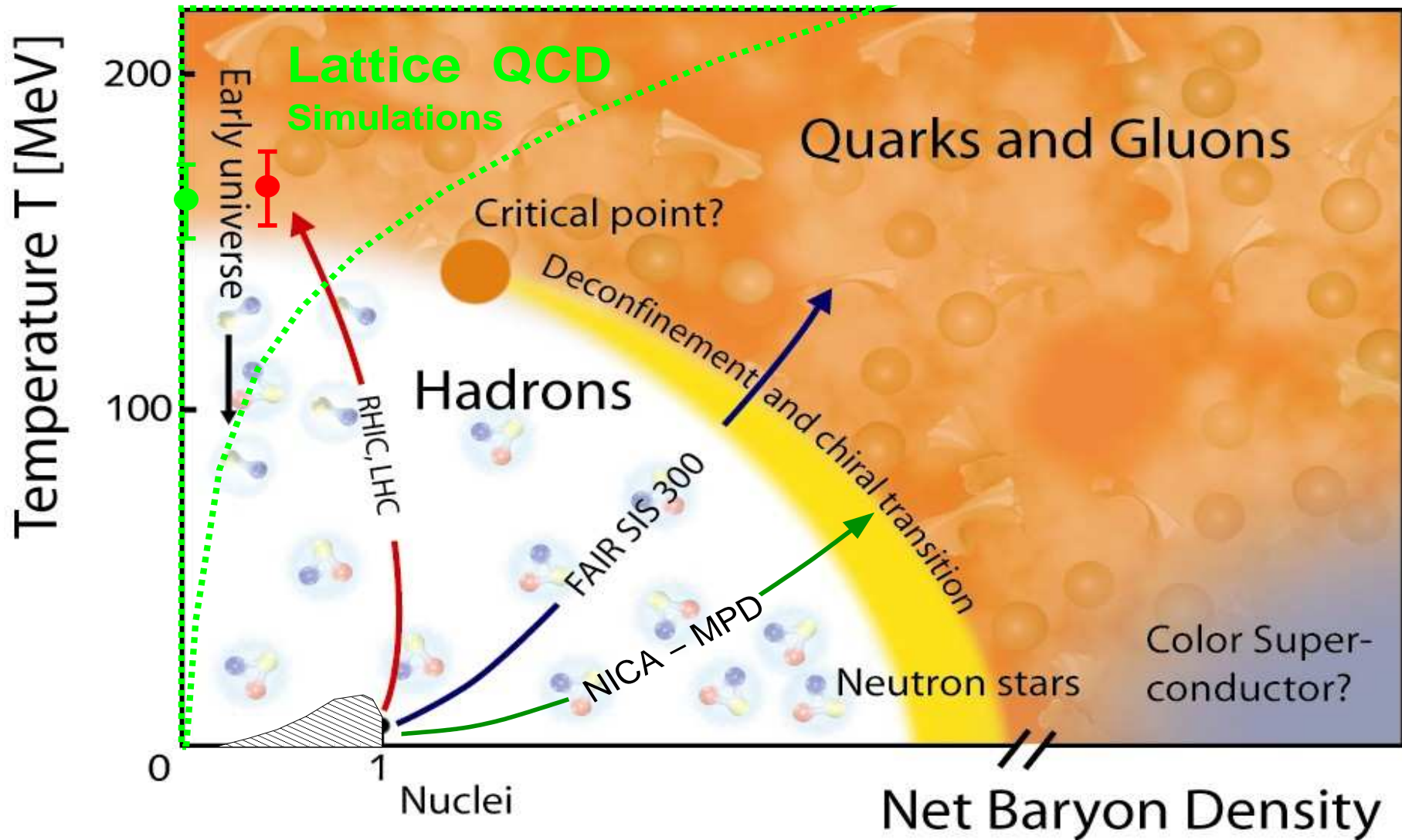
$$\epsilon/T^4 = \frac{\pi^2}{30} N_\pi \sim 1 \quad (\text{ideal pion gas})$$

$$\epsilon/T^4 = \frac{\pi^2}{30} (N_G + \frac{7}{8} N_Q) \sim 15.6 \quad (\text{quarks and gluons})$$

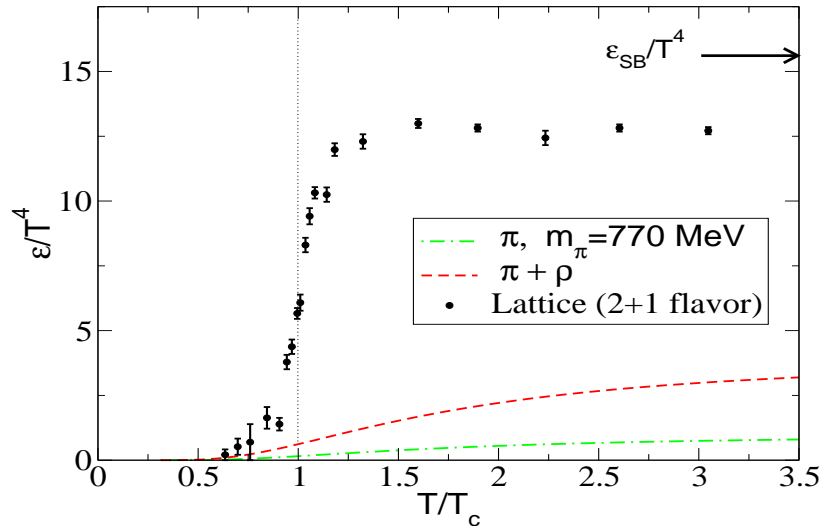
PHASE DIAGRAM OF QCD: LATTICE SIMULATIONS



PHASE DIAGRAM OF QCD: LATTICE SIMULATIONS



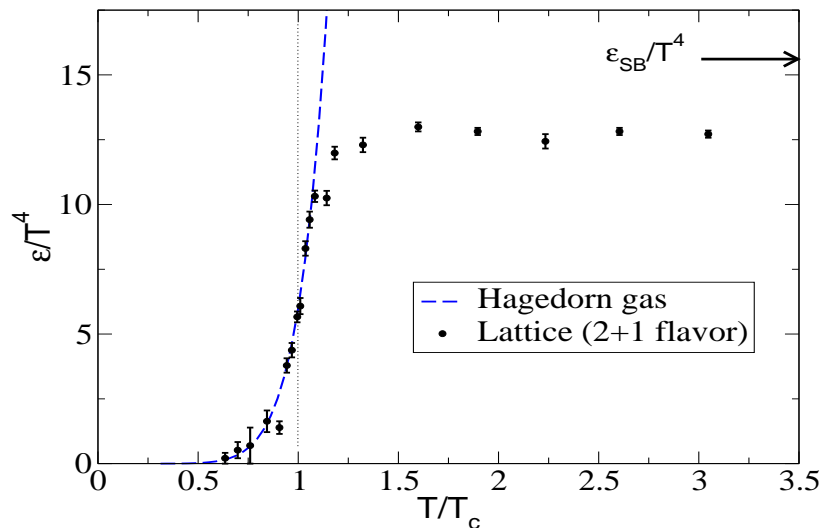
LATTICE QCD EoS VS. RESONANCE GAS



Ideal hadron gas mixture ...

$$\varepsilon(T) = \sum_{i=\pi,\rho,\dots} g_i \int \frac{d^3p}{(2\pi)^3} \frac{\sqrt{p^2 + m_i^2}}{\exp(\sqrt{p^2 + m_i^2}/T) + \delta_i}$$

missing degrees of freedom below and above T_c



Resonance gas ...

Karsch, Redlich, Tawfik, Eur.Phys.J. C29, 549 (2003)

$$\varepsilon(T) = \sum_{i=\pi,\rho,\dots} \varepsilon_i(T) + \sum_{r=M,B} g_r \int dm \rho(m) \int \frac{d^3p}{(2\pi)^3} \frac{\sqrt{p^2 + m^2}}{\exp(\sqrt{p^2 + m^2}/T) + \delta_r}$$

$\rho(m) \sim m^\beta \exp(m/T_H)$... Hagedorn mass spektrum

too many degrees of freedom above T_c

LATTICE QCD EoS AND MOTT-HAGEDORN GAS

$$\varepsilon_R(T, \{\mu_j\}) = \sum_{i=\pi, K, \dots} \varepsilon_i(T, \{\mu_i\}) + \sum_{r=M, B} g_r \int_{m_r} dm \int ds \rho(m) A(s, m; T) \int \frac{d^3 p}{(2\pi)^3} \frac{\sqrt{p^2 + s}}{\exp\left(\frac{\sqrt{p^2 + s} - \mu_r}{T}\right) + \delta_r}$$

Hagedorn mass spectrum: $\rho(m)$

Spectral function for heavy resonances:

$$A(s, m; T) = N_s \frac{m \Gamma(T)}{(s - m^2)^2 + m^2 \Gamma^2(T)}$$

Ansatz with **Mott effect** at $T = T_H = 180$ MeV:

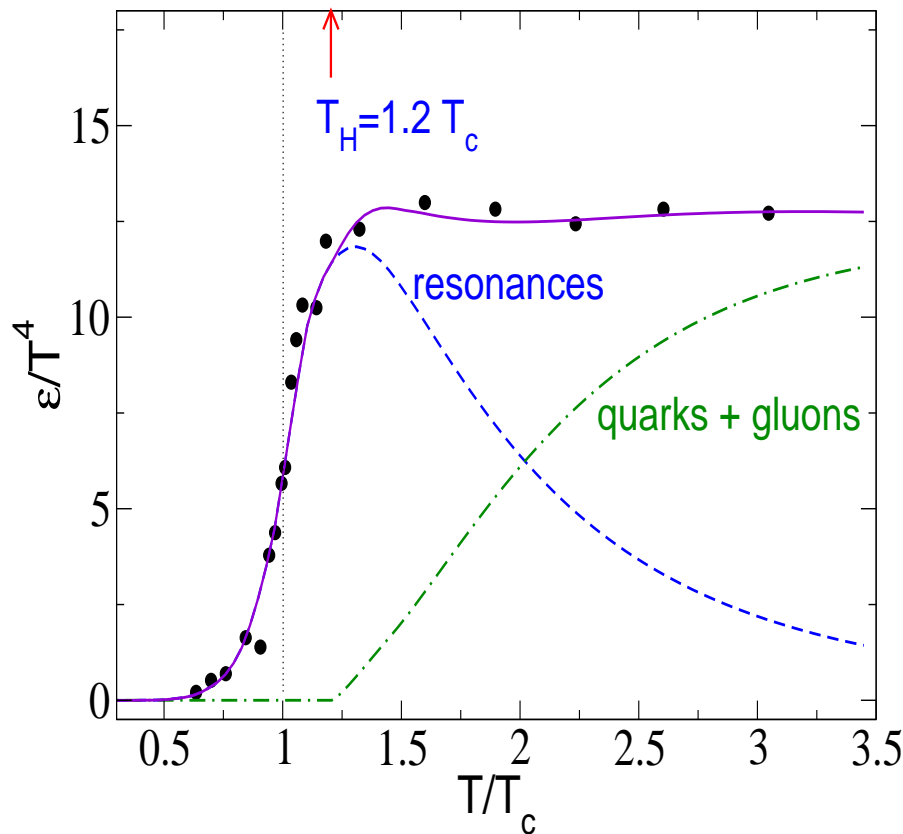
$$\Gamma(T) = B \Theta(T - T_H) \left(\frac{m}{T_H}\right)^{2.5} \left(\frac{T}{T_H}\right)^6 \exp\left(\frac{m}{T_H}\right)$$

No width below T_H : Hagedorn resonance gas
Apparent phase transition at $T_c \sim 150$ MeV

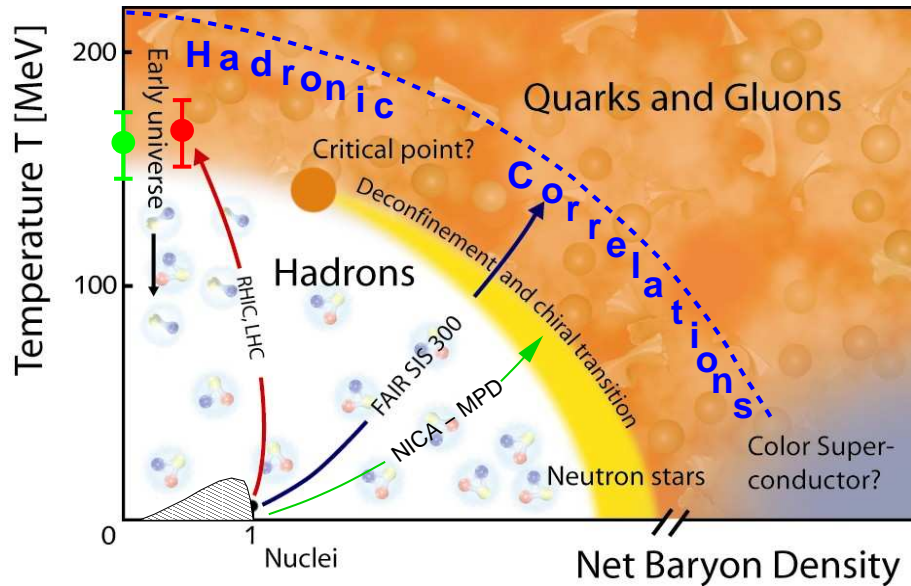
Blaschke & Bugaev, Fizika B13, 491 (2004)

Prog. Part. Nucl. Phys. 53, 197 (2004)

Blaschke & Yudichev, in preparation



HADRONIC CORRELATIONS ABOVE T_c : LATTICE QCD



Hadron correlators $G_H \Rightarrow$ **spectral densities** $\rho_H(\omega, T)$

$$G_H(\tau, T) = \int_0^\infty d\omega \rho_H(\omega, T) \frac{\cosh(\omega(\tau - T/2))}{\sinh(\omega/2T)}$$

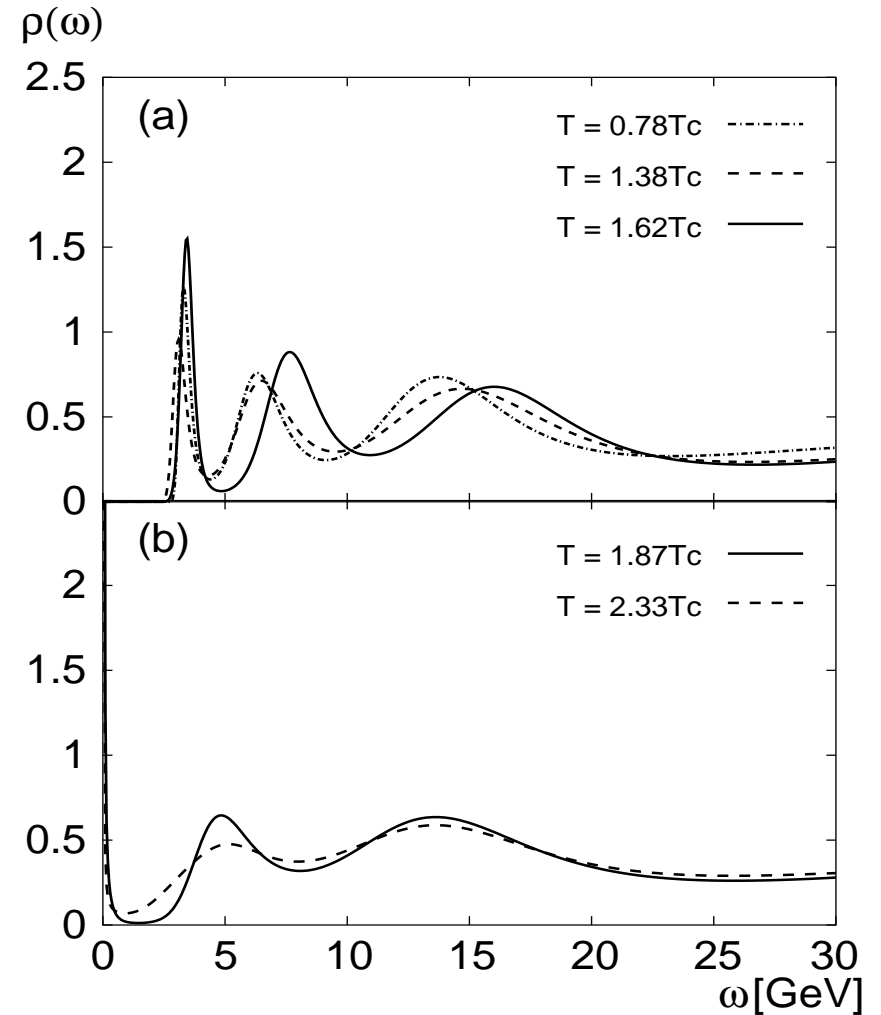
Maximum entropy method

Karsch et al. PLB 530 (2002) 147

Result:

Correlations persist above T_c !

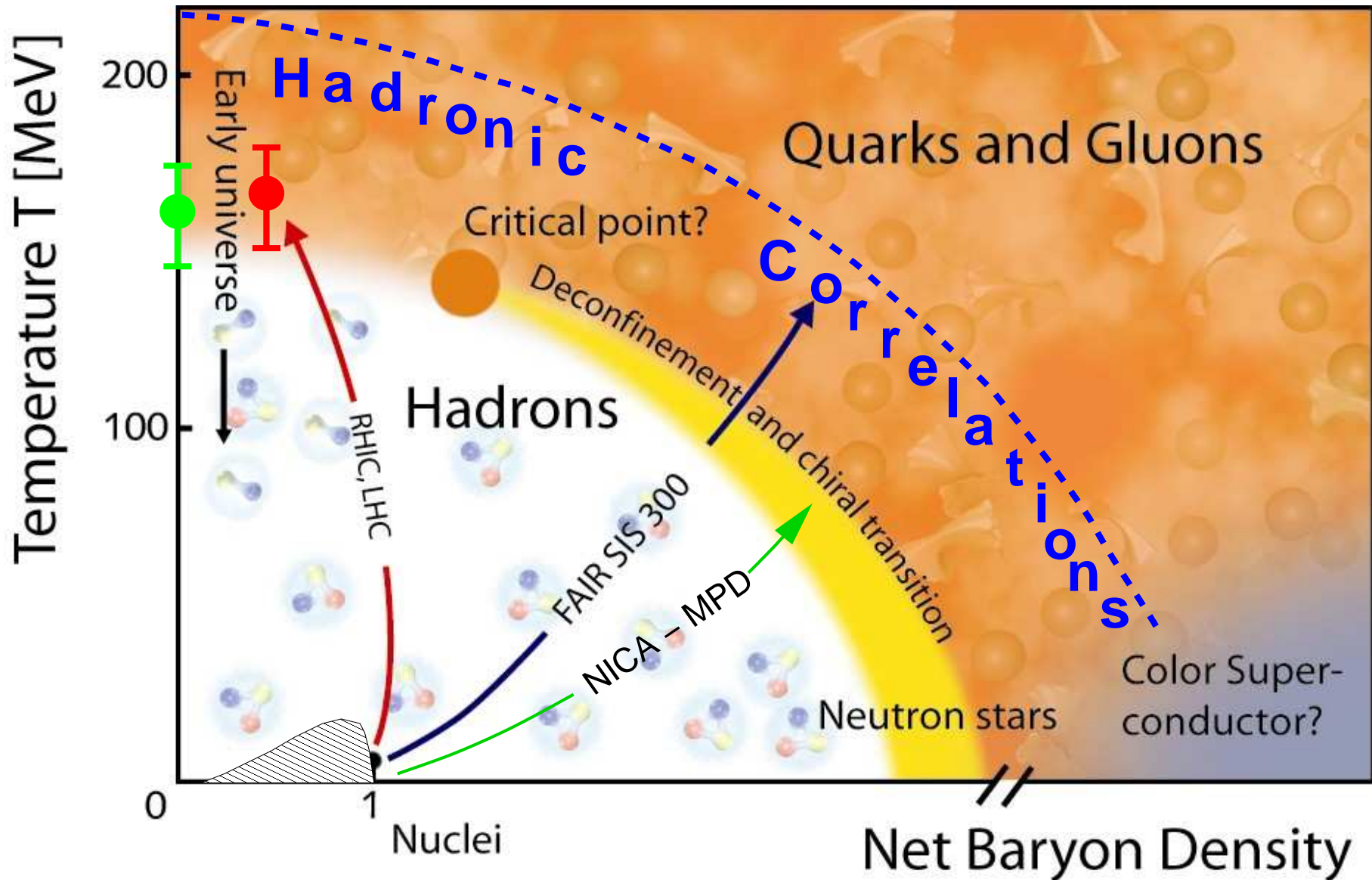
Karsch et al. NPA 715 (2003)



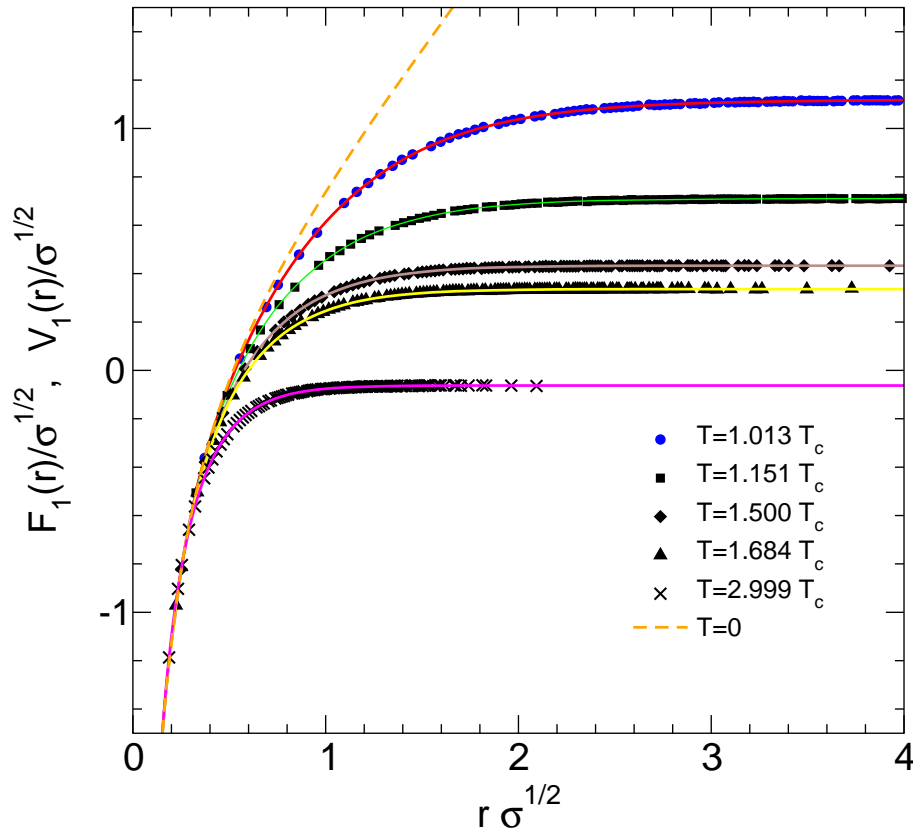
J/ψ and η_c survive up to $T \sim 1.6T_c$

Asakawa, Hatsuda; PRL 92 (2004) 012001

HADRONIC CORRELATIONS IN THE PHASE DIAGRAM OF QCD



HEAVY QUARK POTENTIAL FROM LATTICE QCD



Blaschke, Kaczmarek, Laermann, Yudichev,
EPJC 43, 81 (2005); [hep-ph/0505053]

Color-singlet free energy F_1 in quenched QCD

$$\langle \text{Tr}[L(0)L^\dagger(r)] \rangle = \exp[-F_1(r)/T]$$

Long- and short- range parts

$$F_1(r, T) = F_{1,\text{long}}(r, T) + V_{1,\text{short}}(r)e^{-(\mu(T)r)^2}$$

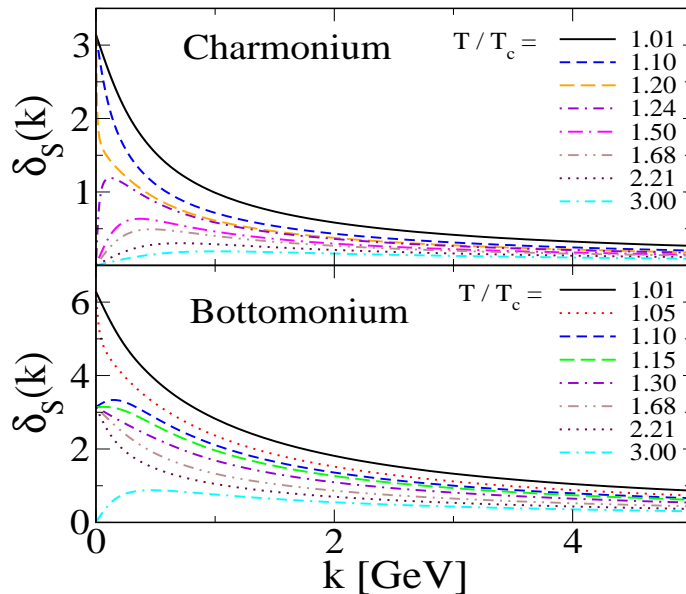
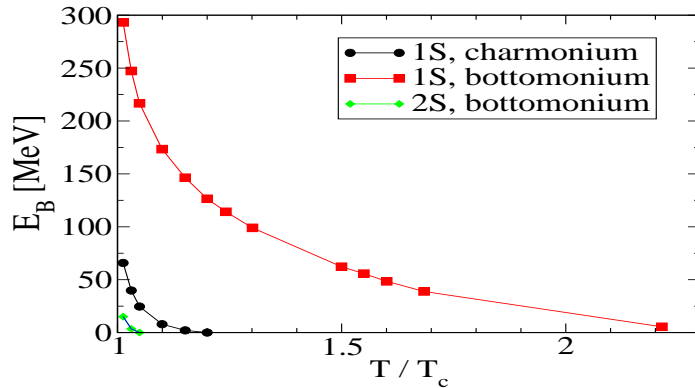
$$F_{1,\text{long}}(r, T) = -\frac{q^2(T)}{2^{3/4}\Gamma(3/4)}\sqrt{\frac{r}{\mu(T)}}K_{1/4}\left[(\mu(T)r)^2\right] + q^2(T)\frac{\Gamma(1/4)}{2^{3/2}\Gamma(3/4)\mu(T)},$$

$$V_{1,\text{short}}(r) = -\frac{4}{3}\frac{\alpha(r)}{r},$$

$$\alpha(r) = \frac{4\pi}{11}\left(\frac{1}{\ln(r^2/c^2)} - \frac{r^2}{r^2 - c^2}\right).$$

Quarkonium ($Q\bar{Q}$)	1S	1P ₁	2S
Charmonium ($c\bar{c}$)	J/ ψ (3097)	χ_{c1} (3510)	ψ' (3686)
Bottomonium ($b\bar{b}$)	Υ (9460)	χ_{b1} (9892)	Υ' (10023)

SCHROEDINGER EQN: BOUND & SCATTERING STATES



Quarkonia **bound states** at finite T :

$$[-\nabla^2/m_Q + V_{\text{eff}}(r, T)]\psi(r, T) = E_B(T)\psi(r, T)$$

Binding energy vanishes $E_B(T_{\text{Mott}}) = 0$: **Mott effect**

Scattering states:

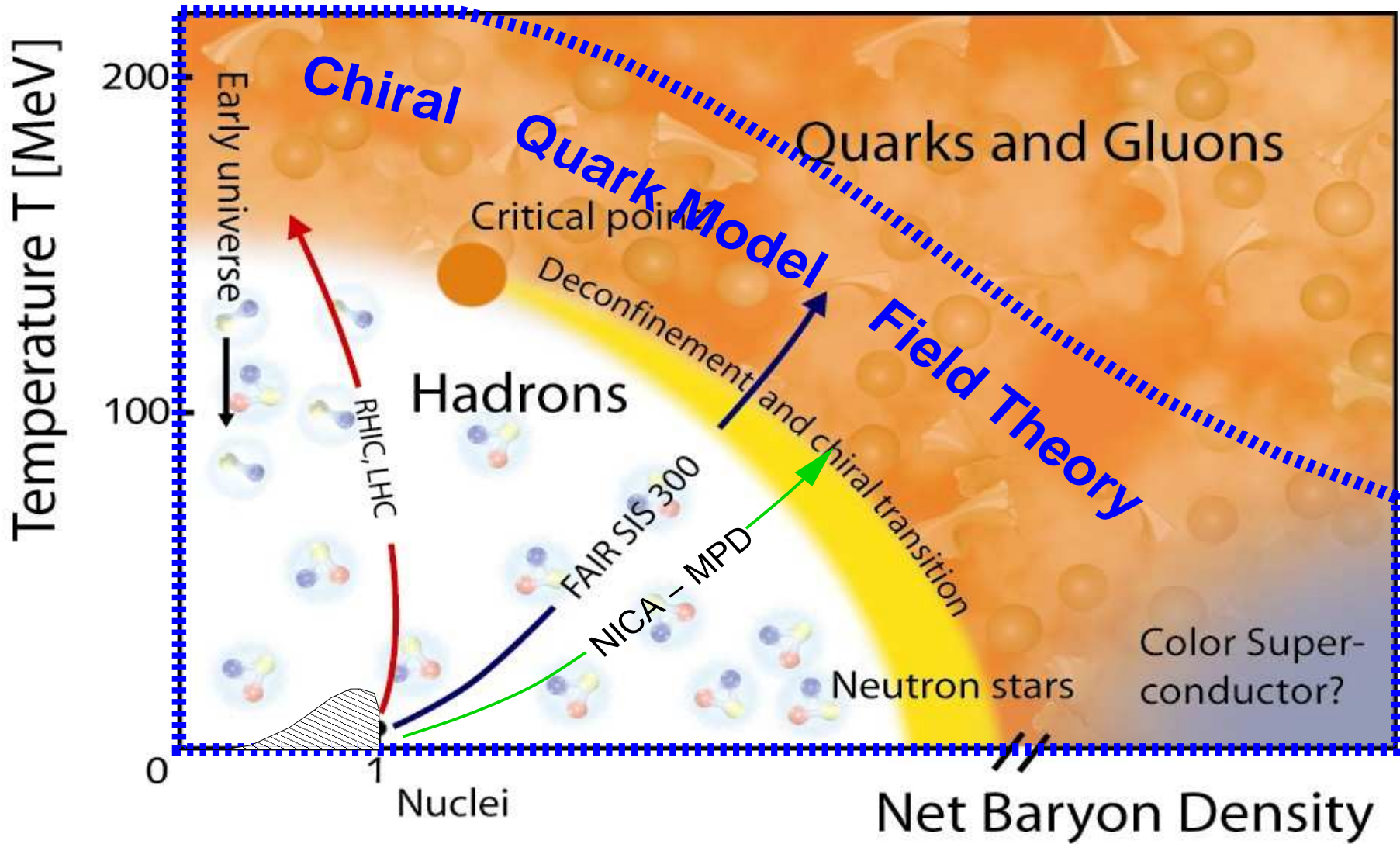
$$\frac{d\delta_S(k, r, T)}{dr} = -\frac{m_Q V_{\text{eff}}}{k} \sin(kr + \delta_S(k, r, T))$$

Levinson theorem:

Phase shift at threshold jumps by π when bound state \rightarrow resonance at $T = T_{\text{Mott}}$

Blaschke, Kaczmarek, Laermann, Yudichev
EPJC 43, 81 (2005); [hep-ph/0505053]

PHASE DIAGRAM OF QCD: CHIRAL MODEL FIELD THEORIES



CHIRAL MODEL FIELD THEORY FOR QUARK MATTER

- Partition function as a Path Integral (imaginary time $\tau = i t$)

$$Z[T, V, \mu] = \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \exp \left\{ - \int^{\beta} d\tau \int_V d^3x [\bar{\psi}(i\gamma^{\mu}\partial_{\mu} - m - \gamma^0\mu)\psi - \mathcal{L}_{\text{int}}] \right\}$$

- Current-current interaction (4-Fermion coupling)

$$\mathcal{L}_{\text{int}} = \sum_{M=\pi,\sigma,\dots} G_M (\bar{\psi}\Gamma_M\psi)^2 + \sum_D G_D (\bar{\psi}^C\Gamma_D\psi)^2$$

- Bosonization (Hubbard-Stratonovich Transformation)

$$Z[T, V, \mu] = \int \mathcal{D}M_M \mathcal{D}\Delta_D^{\dagger} \mathcal{D}\Delta_D \exp \left\{ - \sum_M \frac{M_M^2}{4G_M} - \sum_D \frac{|\Delta_D|^2}{4G_D} + \frac{1}{2} \text{Tr} \ln S^{-1}[\{M_M\}, \{\Delta_D\}] \right\}$$

- Collective (stochastic) fields: Mesons (M_M) and Diquarks (Δ_D)

- Systematic evaluation: **Mean fields** + **Fluctuations**

- Mean-field approximation: **order parameters** for phase transitions (gap equations)
- Lowest order fluctuations: **hadronic correlations** (bound & scattering states)
- Higher order fluctuations: hadron-hadron **interactions**

NJL MODEL FOR NEUTRAL 3-FLAVOR QUARK MATTER

Thermodynamic Potential $\Omega(T, \mu) = -T \ln Z[T, \mu]$

$$\Omega(T, \mu) = \frac{\phi_u^2 + \phi_d^2 + \phi_s^2}{8G_S} + \frac{|\Delta_{ud}|^2 + |\Delta_{us}|^2 + |\Delta_{ds}|^2}{4G_D} - T \sum_n \int \frac{d^3p}{(2\pi)^3} \frac{1}{2} \text{Tr} \ln \left(\frac{1}{T} S^{-1}(i\omega_n, \vec{p}) \right) + \Omega_e - \Omega_0.$$

Inverse Nambu – Gorkov Propagator $S^{-1}(i\omega_n, \vec{p}) = \begin{bmatrix} \gamma_\mu p^\mu - M(\vec{p}) + \mu\gamma^0 & \hat{\Delta}(\vec{p}) \\ \hat{\Delta}^\dagger(\vec{p}) & \gamma_\mu p^\mu - M(\vec{p}) - \mu\gamma^0 \end{bmatrix},$

$$\hat{\Delta}(\vec{p}) = i\gamma_5 \epsilon_{\alpha\beta\gamma} \epsilon_{ijk} \Delta_{k\gamma} g(\vec{p}) \quad ; \quad \Delta_{k\gamma} = 2G_D \langle \bar{q}_{i\alpha} i\gamma_5 \epsilon_{\alpha\beta\gamma} \epsilon_{ijk} g(\vec{q}) q_{j\beta}^C \rangle.$$

Fermion Determinant (Tr ln D = ln det D): $\ln \det[\beta S^{-1}(i\omega_n, \vec{p})] = 2 \sum_{a=1}^{18} \ln \{ \beta^2 [\omega_n^2 + \lambda_a(\vec{p})^2] \} .$

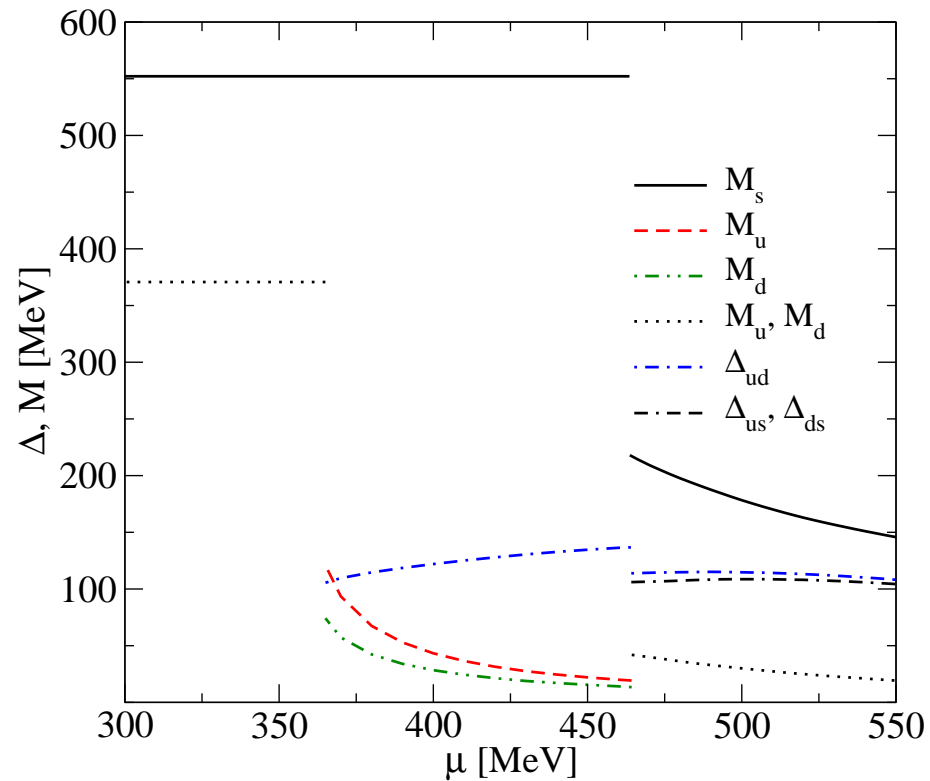
Result for the thermodynamic Potential (Meanfield approximation)

$$\Omega(T, \mu) = \frac{\phi_u^2 + \phi_d^2 + \phi_s^2}{8G_S} + \frac{|\Delta_{ud}|^2 + |\Delta_{us}|^2 + |\Delta_{ds}|^2}{4G_D} - \int \frac{d^3p}{(2\pi)^3} \sum_{a=1}^{18} \left[\lambda_a + 2T \ln \left(1 + e^{-\lambda_a/T} \right) \right] + \Omega_e - \Omega_0.$$

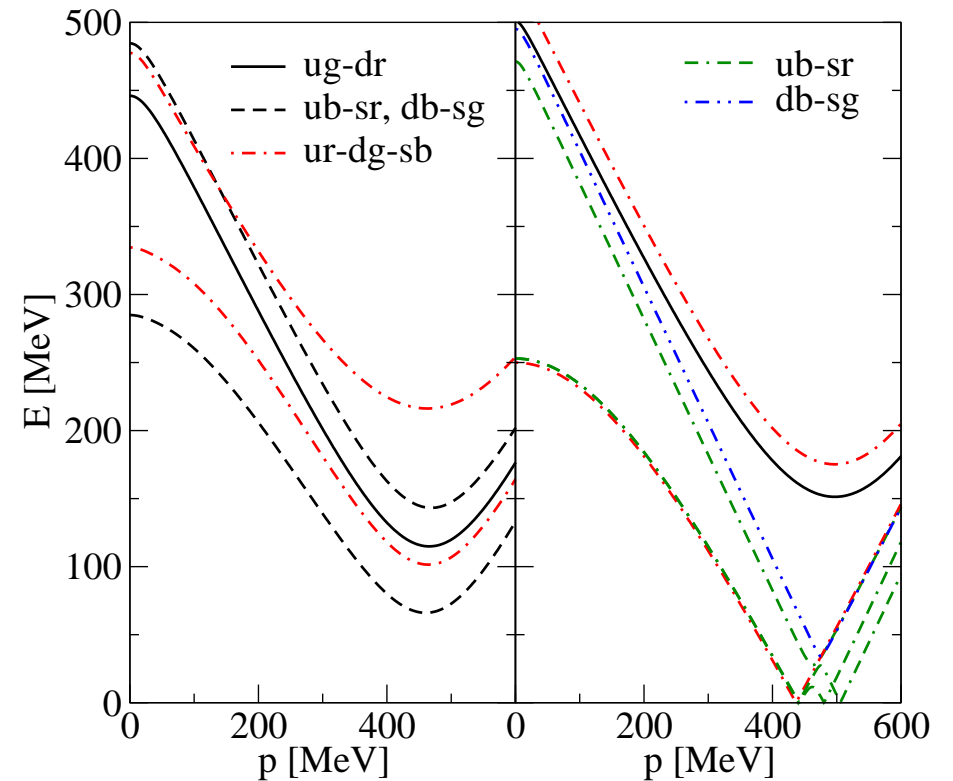
Color and electric charge neutrality constraints: $n_Q = n_8 = n_3 = 0, n_i = -\partial\Omega/\partial\mu_i = 0,$
Equations of state: $P = -\Omega,$ etc.

ORDER PARAMETERS: MASSES AND DIQUARK GAPS

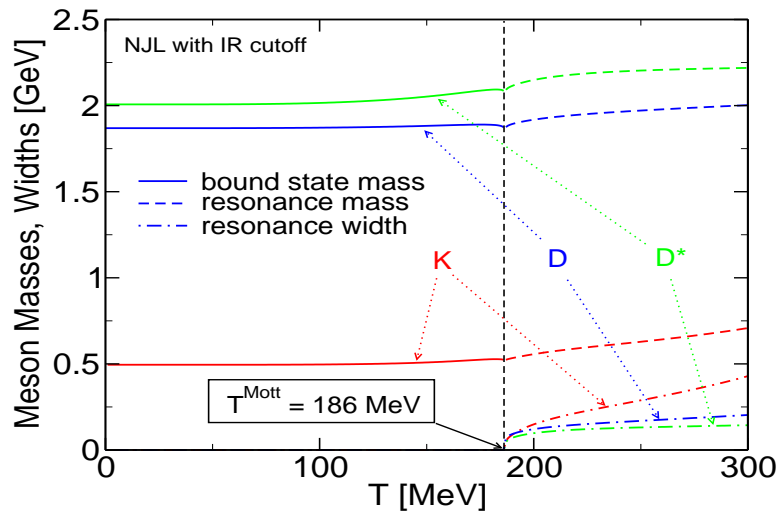
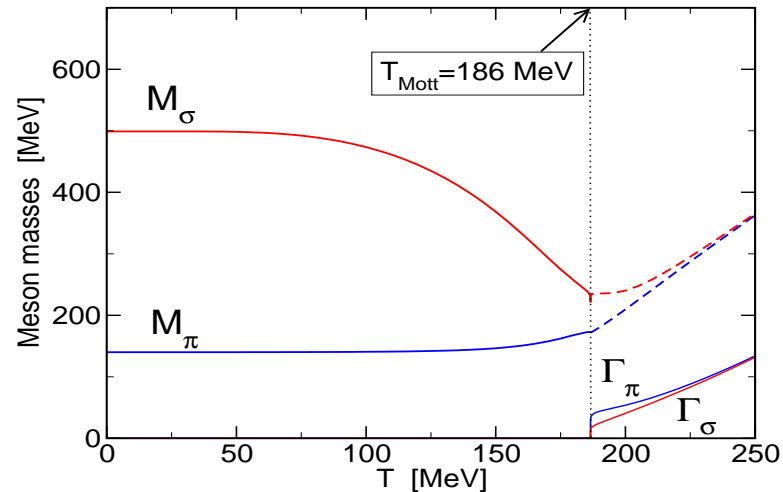
Masses (M) and Diquark gaps (Δ) as a function of the chemical potential at $T = 0$



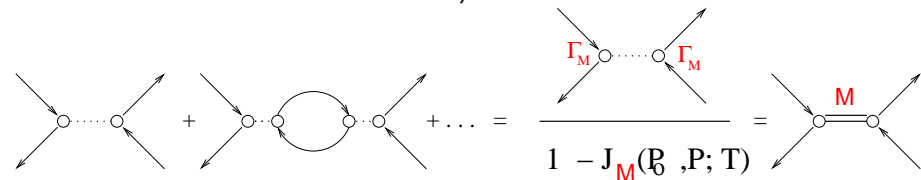
Left: Gap in excitation spectrum ($T = 0$)
Right: 'Gapless' excitations ($T = 60$ MeV)



MOTT EFFECT: NJL MODEL PRIMER



RPA-type resummation of quark-antiquark scattering in the mesonic channel M ,



defines Meson propagator

$$D_M(P_0, P; T) \sim [1 - J_M(P_0, P; T)]^{-1},$$

by the complex polarization function J_M
 → Breit-Wigner type spectral function

$$\begin{aligned} \mathcal{A}_M(P_0, P; T) &= \frac{1}{\pi} \text{Im} D_M(P_0, P; T) \\ &\sim \frac{1}{\pi} \frac{\Gamma_M(T) M_M(T)}{(s - M_M^2(T))^2 + \Gamma_M^2(T) M_M^2(T)} \end{aligned}$$

For $T < T_{\text{Mott}}$: $\Gamma \rightarrow 0$, i.e. bound state

$$\mathcal{A}_M(P_0, P; T) = \delta(s - M_M^2(T))$$

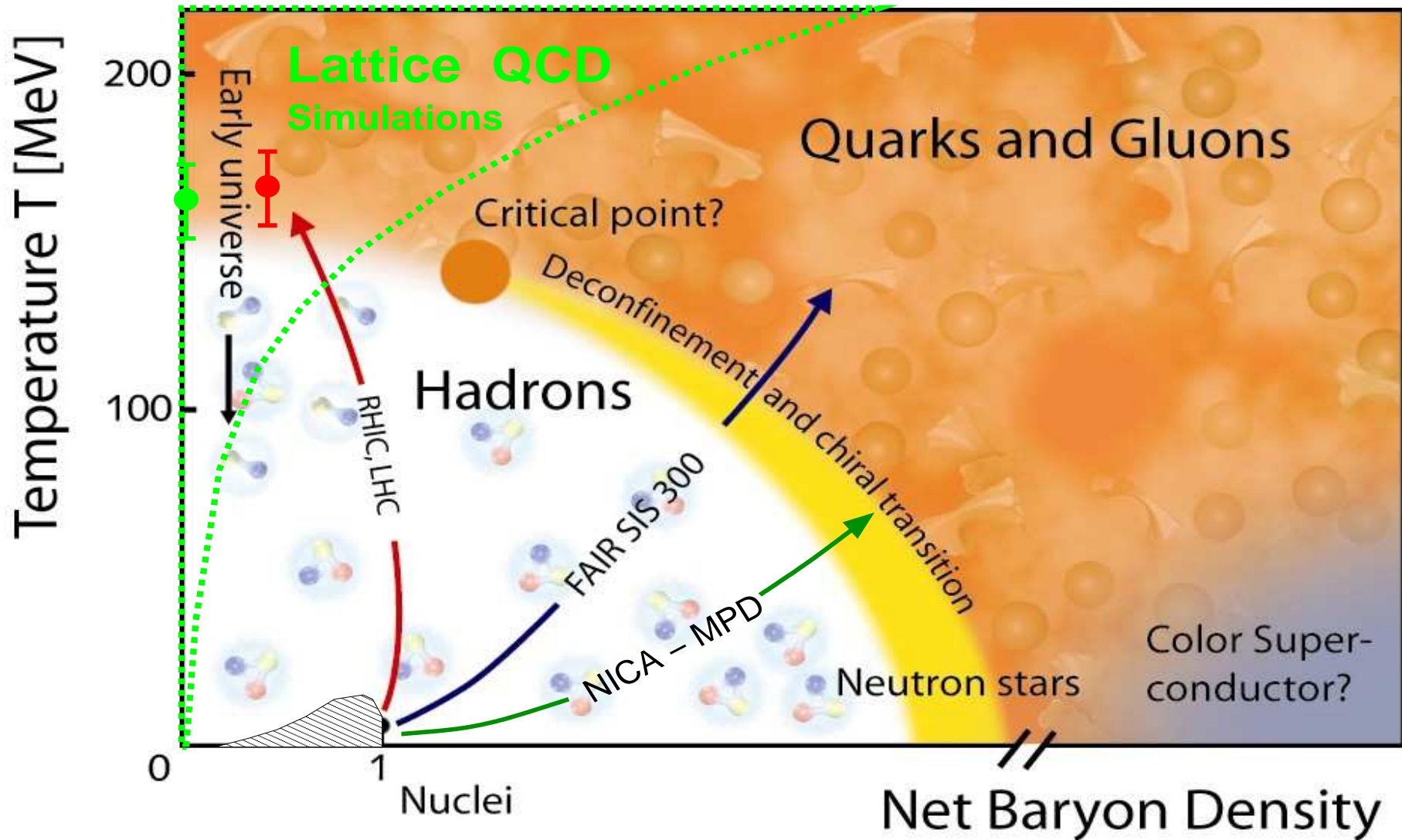
Light meson sector:

Blaschke, Burau, Volkov, Yudichev: EPJA 11 (2001) 319

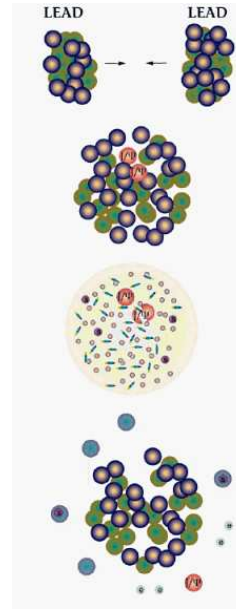
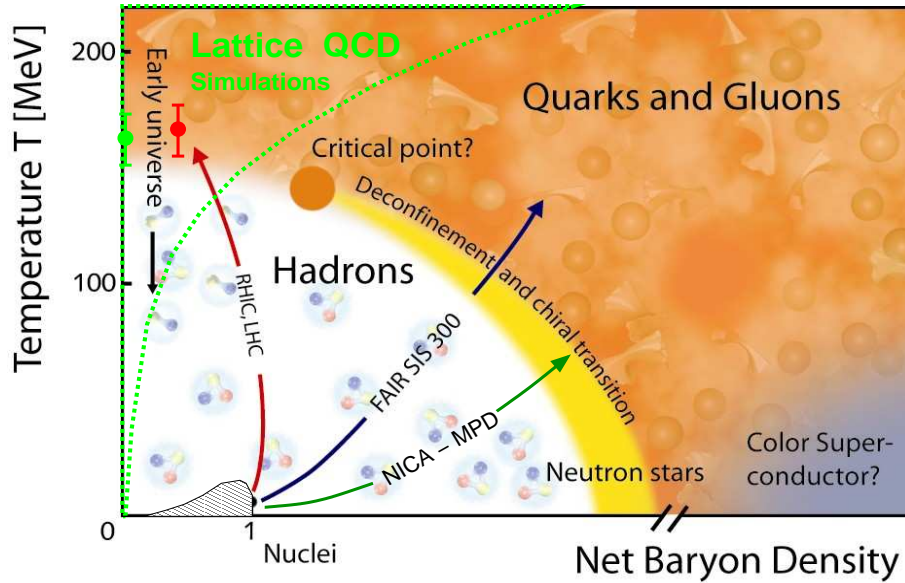
Charm meson sector:

Blaschke, Burau, Kalinovsky, Yudichev, Prog. Theor. Phys. Suppl. 149 (2003) 182

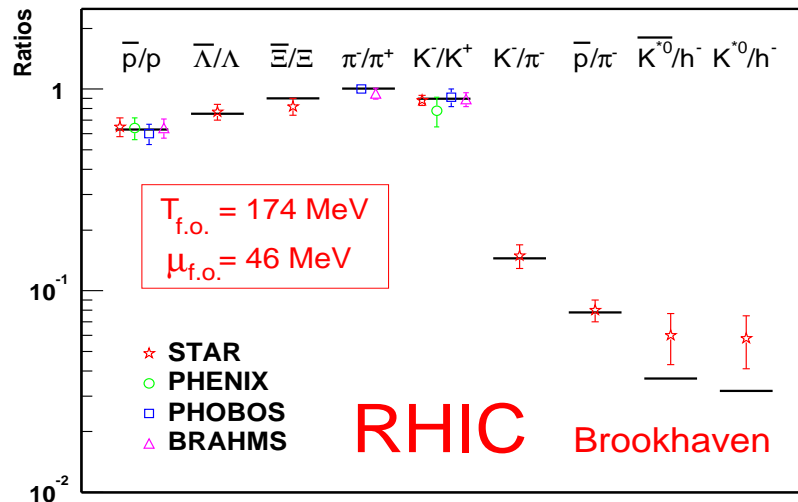
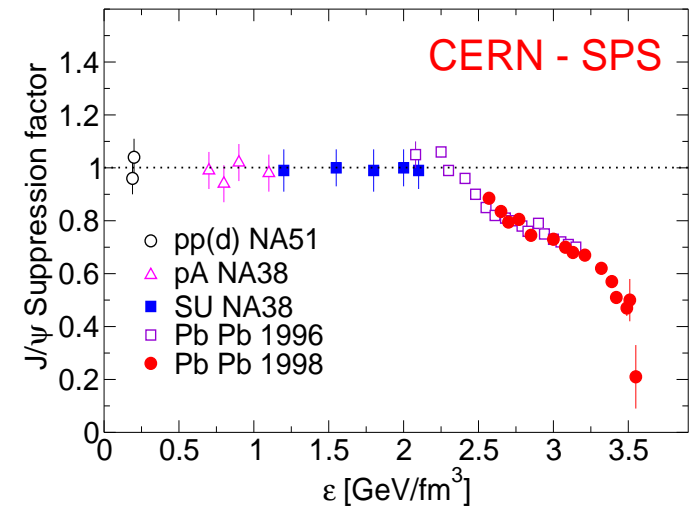
PHASE DIAGRAM OF QCD: HEAVY-ION COLLISIONS



PHASEDIAGRAM OF QCD: LATTICE VS. HEAVY-ION COLLISIONS



QGP Signal: Anomalous J/ψ suppression



Statistical model describes composition of hadron yields in Heavy-Ion Collisions with few freeze-out parameters.

$$\ln Z[T, V, \{\mu\}] = \pm V \sum_i \frac{g_i}{2\pi^2} \int_0^\infty dp p^2 \ln[1 \pm \lambda_i \exp(-\beta \epsilon_i(p))]$$

$$\lambda_i(T, \{\mu\}) = \exp[\beta(\mu_B B_i + \mu_S S_i + \mu_Q Q_i)]$$

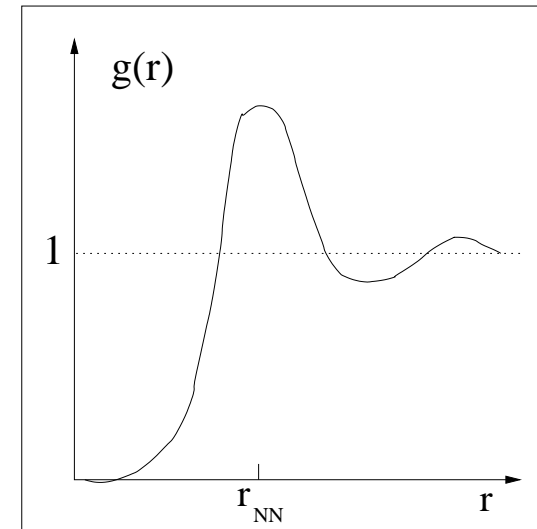
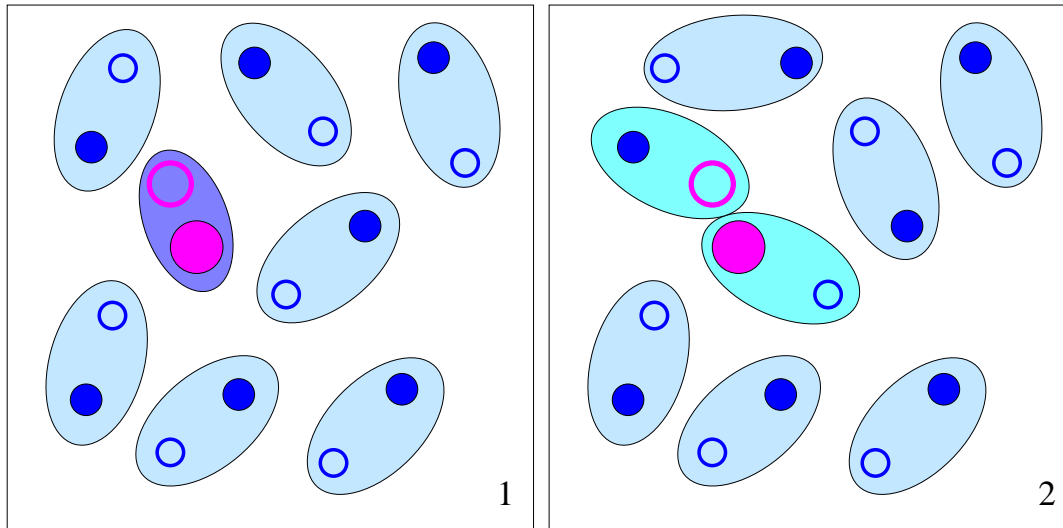
Braun-Munzinger, Redlich, Stachel, in *QGP III* (2003)

A SNAPSHOT OF THE SQGP

The Picture: String-flip (Rearrangement)



Pair correlation

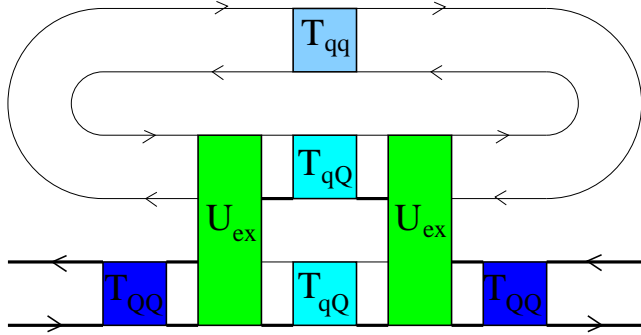


Horowitz et al. PRD (1985), D.B. et al. PLB (1985),
Röpke, Blaschke, Schulz, PRD (1986)

Thoma, Quark Matter '05;
[hep-ph/0509154]

- Strong correlations present: hadronic spectral functions above T_c (lattice QCD)
- Finite width due to rearrangement collisions (higher order correlations)
- Liquid-like pair correlation function (nearest neighbor peak)

QUANTUM KINETIC APPROACH TO J/ψ BREAKUP



Inverse lifetime for Charmonium states

$$\tau^{-1}(p) = \Gamma(p) = \Sigma^>(p) \mp \Sigma^<(p)$$

$$\Sigma^>(p, \omega) = \int_{p'} \int_{p_1} \int_{p_2} (2\pi)^4 \delta_{p,p';p_1,p_2} |\mathcal{M}|^2 G_\pi^<(p') G_{D_1}^>(p_1) G_{D_2}^>(p_2)$$

$$G_h^>(p) = [1 \pm f_h(p)] A_h(p) \text{ and } G_h^<(p) = f_h(p) A_h(p)$$

$$\tau^{-1}(p) = \int \frac{d^3\mathbf{p}'}{(2\pi)^3} \int ds' f_\pi(\mathbf{p}', s') A_\pi(s') v_{\text{rel}} \sigma^*(s)$$

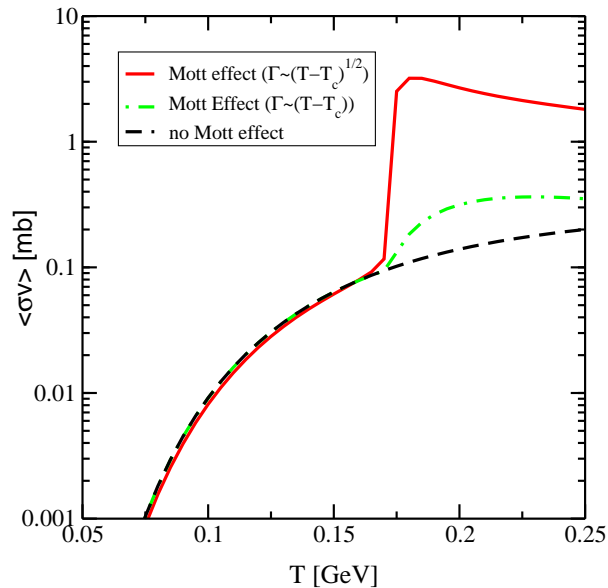
In-medium breakup cross section

$$\sigma^*(s) = \int ds_1 ds_2 A_{D_1}(s_1) A_{D_2}(s_2) \sigma(s; s_1, s_2)$$

Medium effects in **spectral functions** A_h and $\sigma(s; s_1, s_2)$

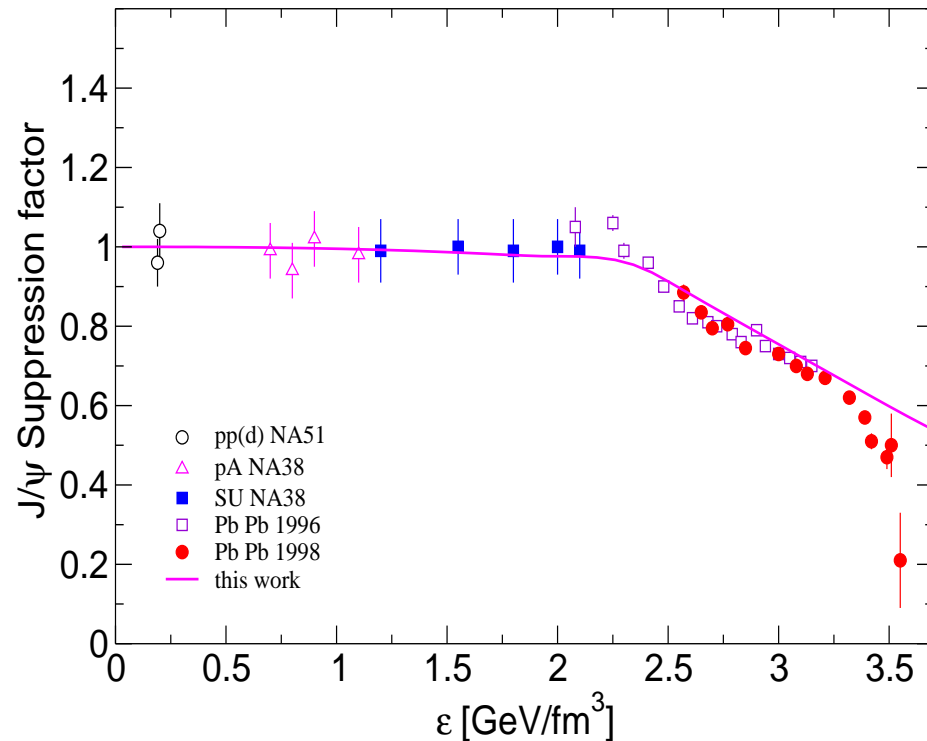
$$A_h(s) = \frac{1}{\pi} \frac{\Gamma_h(T) M_h(T)}{(s - M_h^2(T))^2 + \Gamma_h^2(T) M_h^2(T)} \longrightarrow \delta(s - M_h^2)$$

resonance \leftarrow **Mott-effect** \leftarrow bound state



Blaschke et al., Heavy Ion Phys. 18 (2003) 49

“ANOMALOUS” J/ψ SUPPRESSION IN MOTT-HAGEDORN GAS



Survival probability for J/ψ

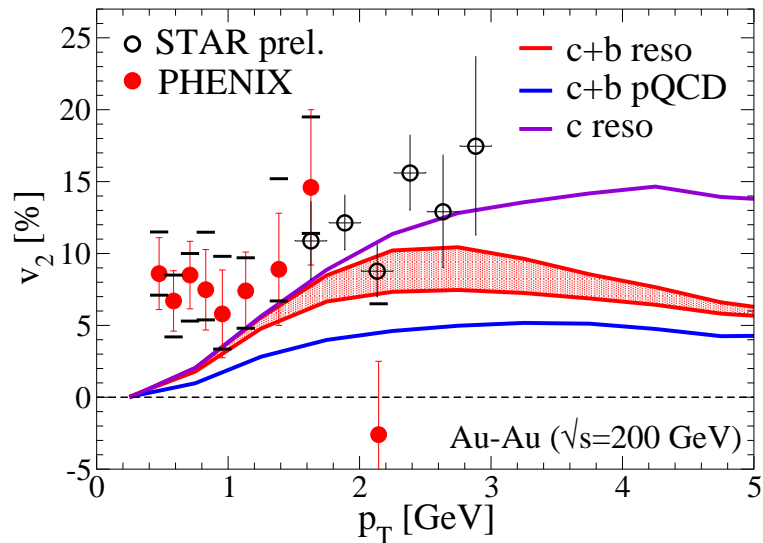
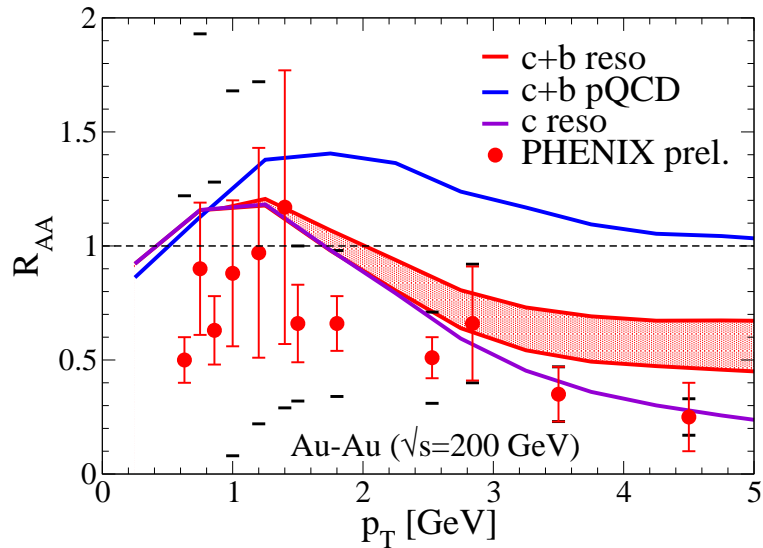
$$S(E_T)/S_N(E_T) = \exp \left[- \int_{t_0}^{t_f} dt \tau^{-1}(n(t)) \right]$$

Threshold: Mott effect for hadrons

Blaschke and Bugaev, Prog. Part. Nucl. Phys. 53 (2004) 197

In progress: full kinetics with gain processes (D-fusion), HIC simulation

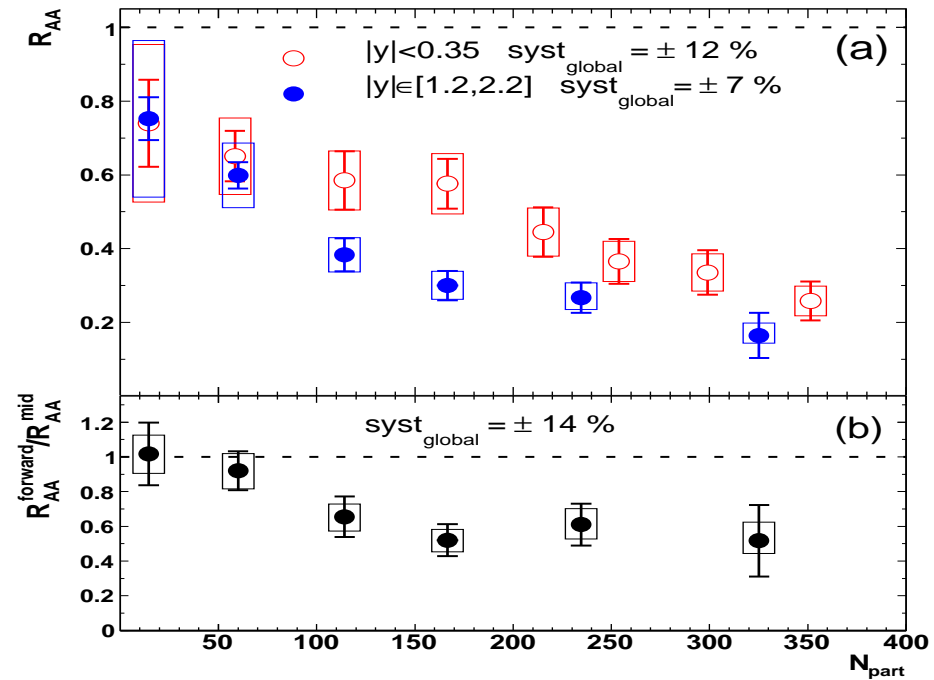
CHARM AND CHARMONIUM PRODUCTION @ RHIC



Nuclear modification factor R_{AA} and elliptic flow v_2 of semileptonic D^- and B^- meson decay-electrons in $b = 7$ fm Au-Au ($\sqrt{s} = 200$ GeV) collisions at RHIC



J/ψ suppression in forward stronger than in central rapidity: signal for charmonium regeneration? ↓



Hees, Greco, Rapp, PRC 73, 034913 (2006)

Adare et al. (PHENIX Collaboration); nucl-ex/0611020

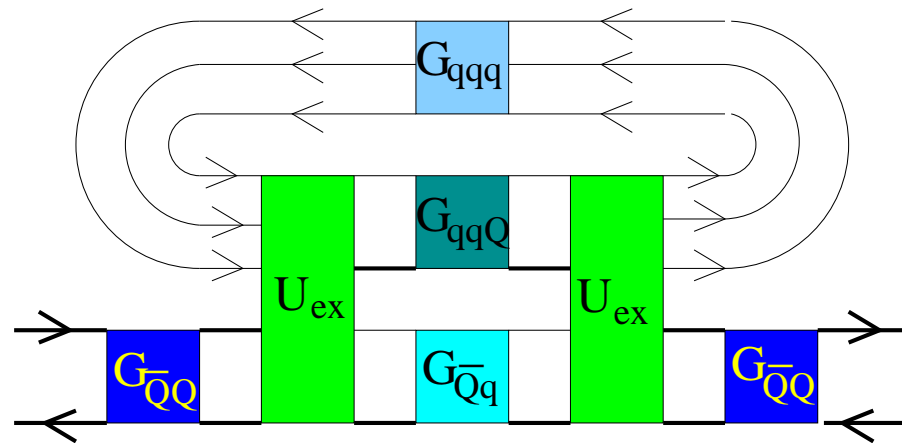
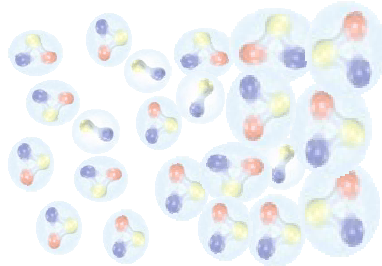
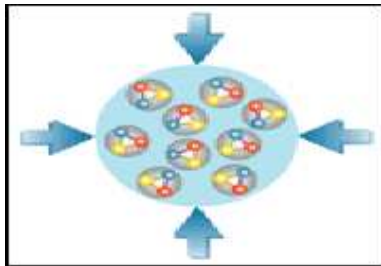
CHARM AND CHARMONIUM PRODUCTION @ FAIR-CBM



J/ψ dissociation process in dense baryonic matter at FAIR-CBM: spectral functions for open charm hadrons (D-meson, Λ_c) are essential inputs!



D-meson spectral function in cold dense nuclear matter from a G-matrix approach ↓

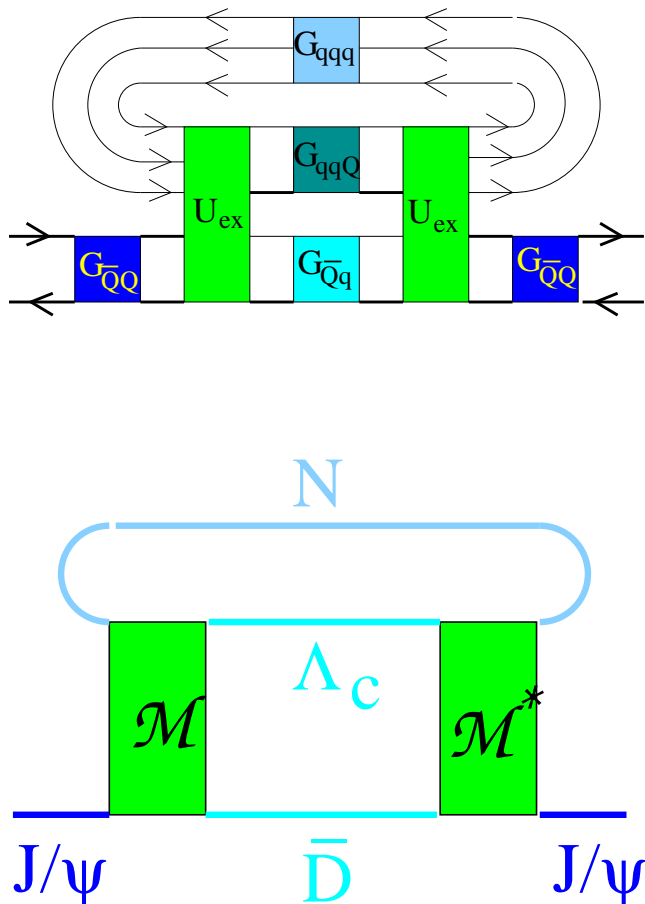


QUANTUM KINETICS OF J/ψ DISSOCIATION @ CBM ($\mu_B \neq 0$)

Charmonium lifetime in a dense nuclear medium ($f_D \approx 0$)

$$\tau^{-1}(p) = \int_{p'} \int_{p_2} \int_{p_2} (2\pi)^4 \delta_{p,p';p_1,p_2} |\mathcal{M}|^2 f_N(p') A_N(p') A_\Lambda(p_1) A_D(p_2)$$

Medium effects in hadronic **spectral functions** A_h and $\sigma(s; s_1, s_2)$
 D-meson spectral function in cold dense nuclear matter from a G-matrix approach \downarrow (N, Λ_c similar)



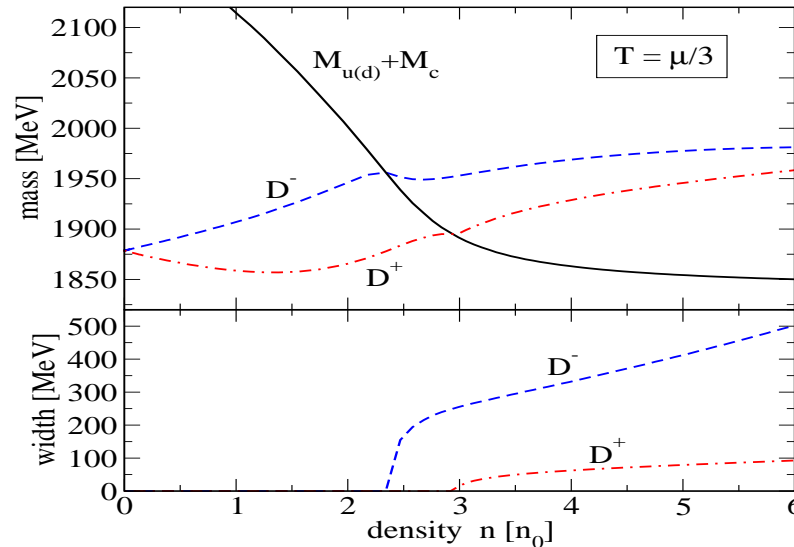
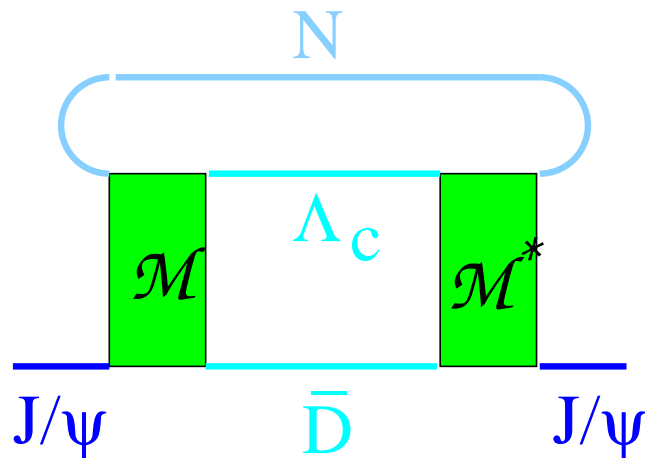
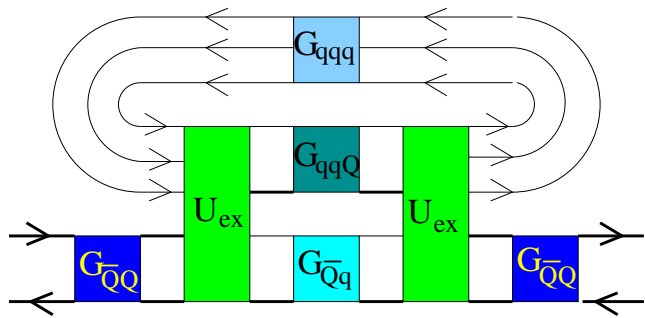
Tolos et al., EPJC (2005); nucl-th/0501151

QUANTUM KINETICS OF J/ψ DISSOCIATION @ CBM ($\mu_B \neq 0$)

Charmonium lifetime in a dense nuclear medium ($f_D \approx 0$)

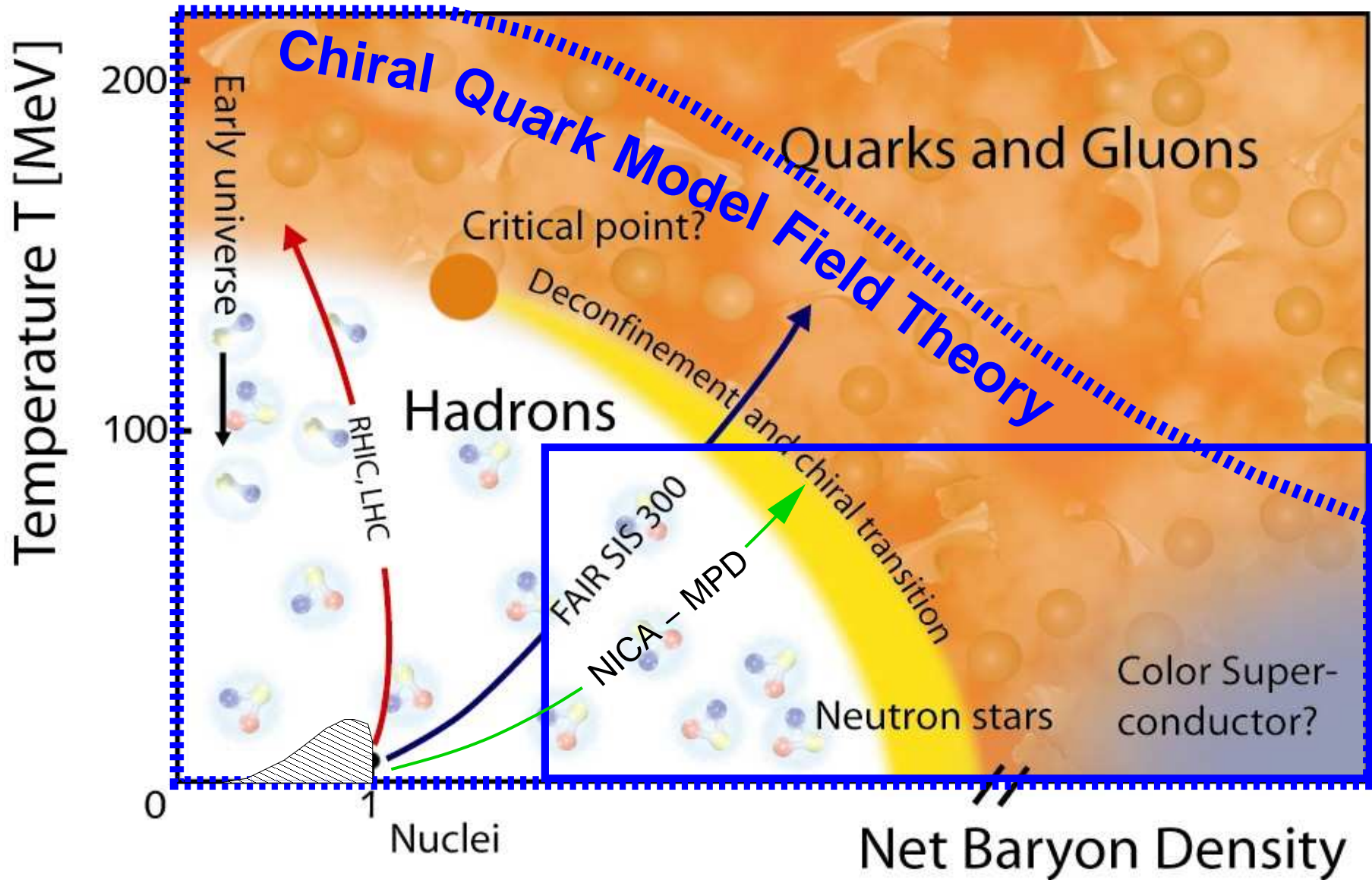
$$\tau^{-1}(p) = \int_{p'} \int_{p_2} \int_{p_2} (2\pi)^4 \delta_{p,p';p_1,p_2} |\mathcal{M}|^2 f_N(p') A_N(p') A_\Lambda(p_1) A_D(p_2)$$

Medium effects in hadronic **spectral functions** A_h and $\sigma(s; s_1, s_2)$
 D-meson spectral function in hot, dense quark matter from a NJL model approach \downarrow (N, Λ_c similar)

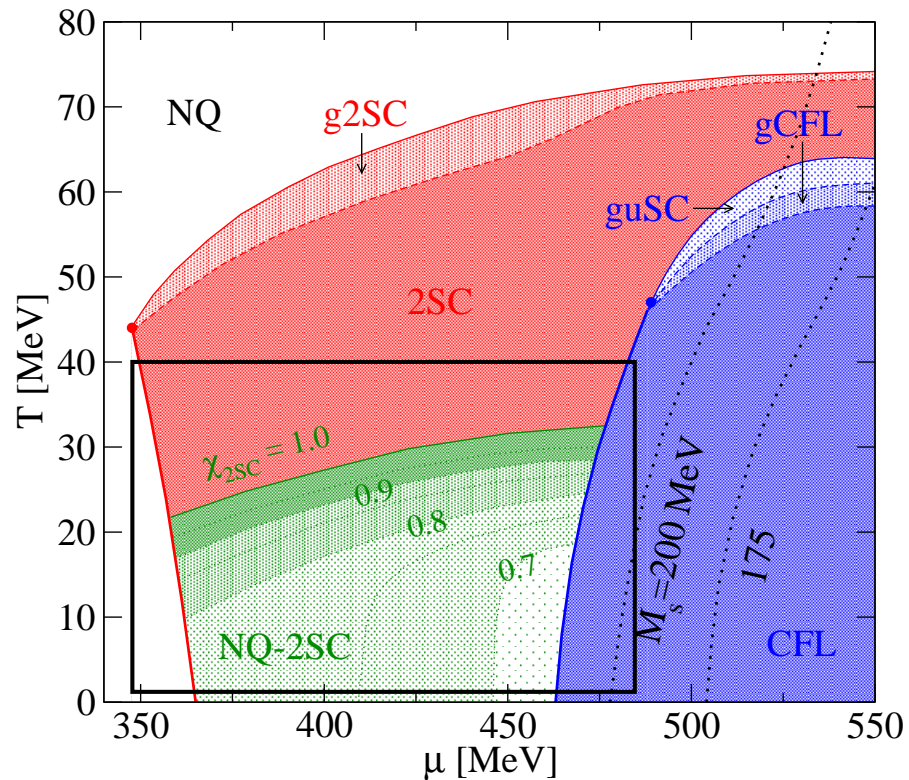


D.B., P. Costa, Yu. Kalinovsky, in preparation

PHASE DIAGRAM OF DEGENERATE QUARK MATTER



QUARK MATTER IN COMPACT STARS



Rüster et al: PRD 72 (2005) 034004
Blaschke et al: PRD 72 (2005) 065020
Abuki, Kunihiro: NPA 768 (2006) 118

The phases are characterized by 3 gaps:

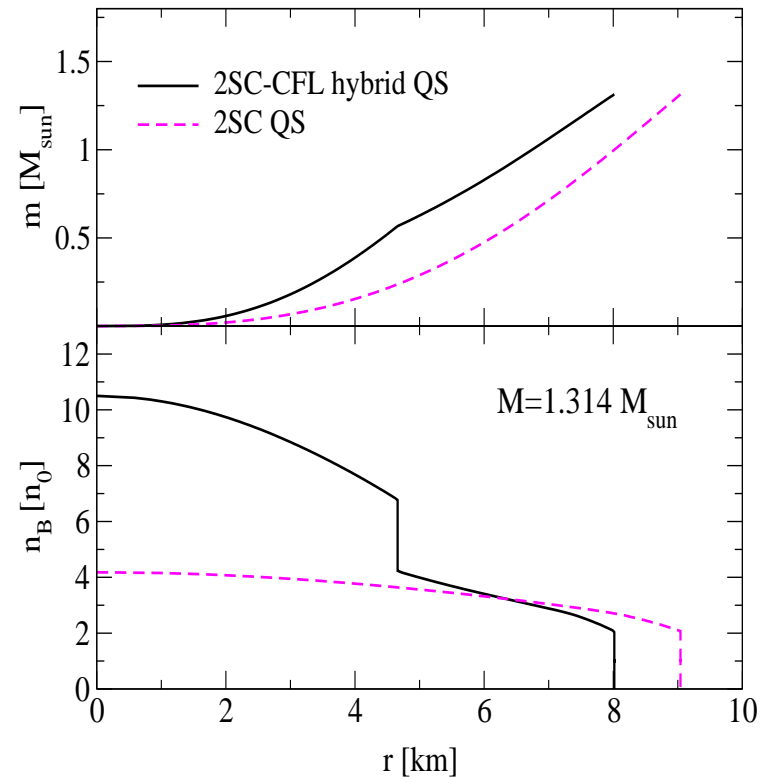
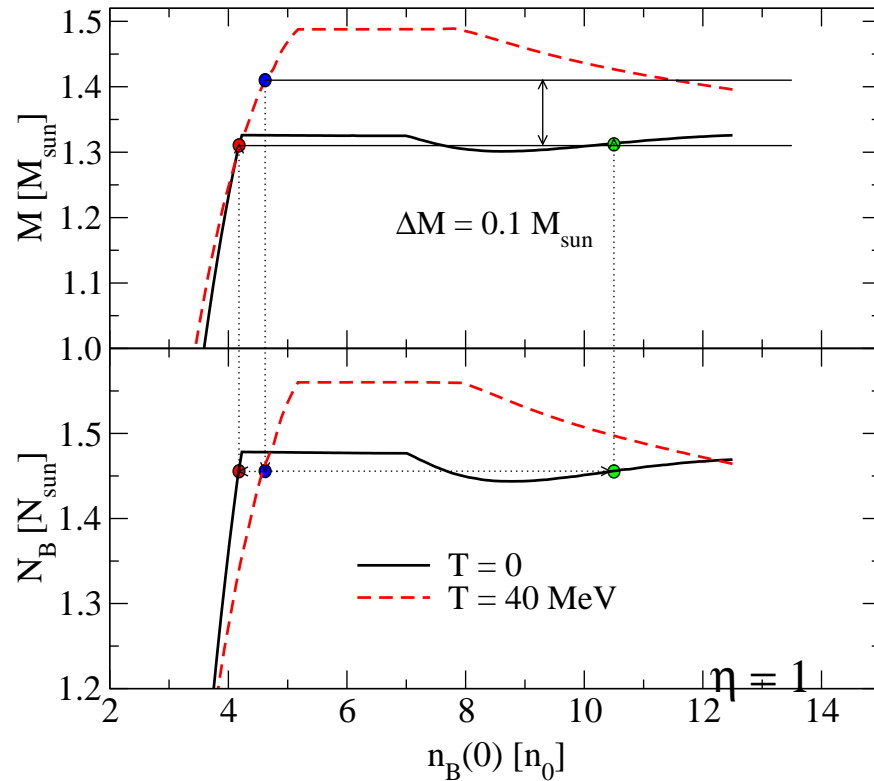
- **NQ:** $\Delta_{ud} = \Delta_{us} = \Delta_{ds} = 0$;
- **NQ-2SC:** $\Delta_{ud} \neq 0, \Delta_{us} = \Delta_{ds} = 0, 0 \leq \chi_{2SC} \leq 1$;
- **2SC:** $\Delta_{ud} \neq 0, \Delta_{us} = \Delta_{ds} = 0$;
- **uSC:** $\Delta_{ud} \neq 0, \Delta_{us} \neq 0, \Delta_{ds} = 0$;
- **CFL:** $\Delta_{ud} \neq 0, \Delta_{ds} \neq 0, \Delta_{us} \neq 0$;

Result:

- Gapless phases only at high T,
- CFL only at high chemical potential,
- At $T \leq 25$ -30 MeV: mixed NQ-2SC phase,
- Critical point $(T_c, \mu_c) = (48 \text{ MeV}, 353 \text{ MeV})$,
- Strong coupling, $\eta = 1$, changes?.

⇒ Zhuang (DM 12, 17)

2SC-CFL TWIN CONFIGURATIONS, ENERGY RELEASE



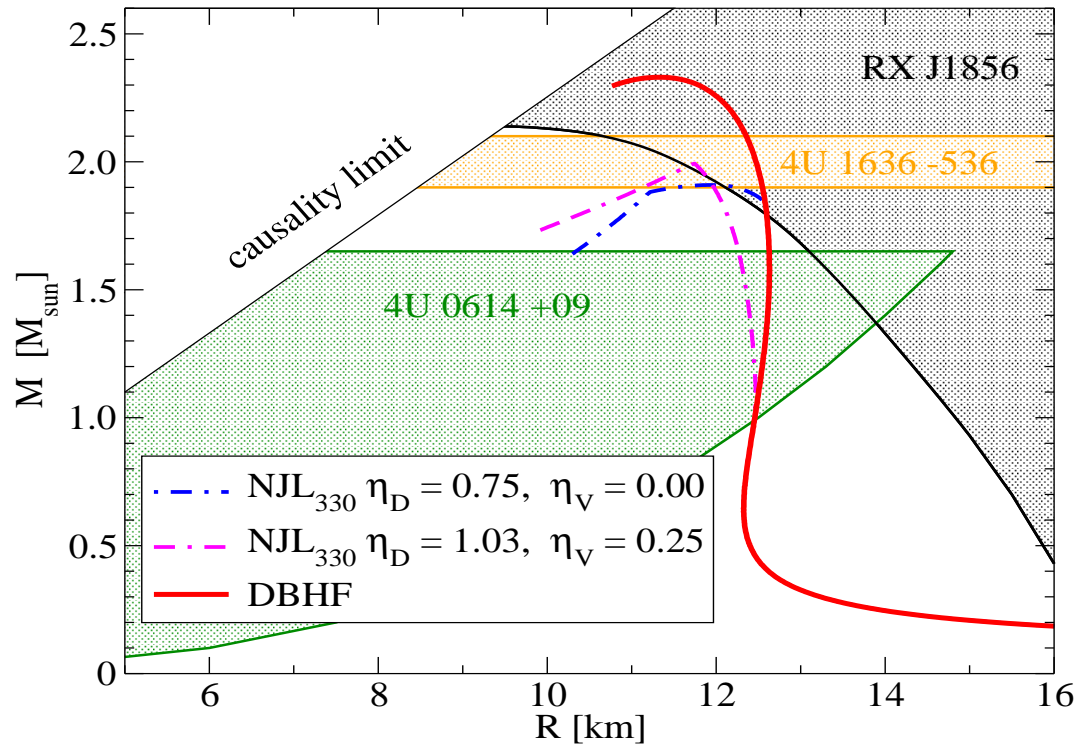
Energy release: $\Delta E \sim 0.1 M_{\odot} c^2 \sim 10^{52}$ erg.

See also: [Aguilera et al: A&A 416, 991 \(2004\)](#), [DB et al: NPA 736, 203 \(2004\)](#)

Caution: CFL core unstable against adding a hadronic shell!

QUARK MATTER IN COMPACT STARS: MASS-RADIUS CONSTRAINT

Solve TOV Eqn. → Hybrid stars fulfill constraint!

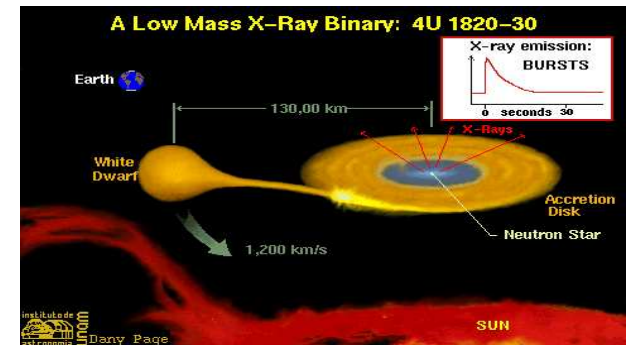


Klähn et al: Constraints on the high-density EoS ...
 PRC 74 (2006); [nucl-th/0602038], [astro-ph/0606524]

- Isolated Neutron star RX J1856:
 M-R constraint from thermal emission

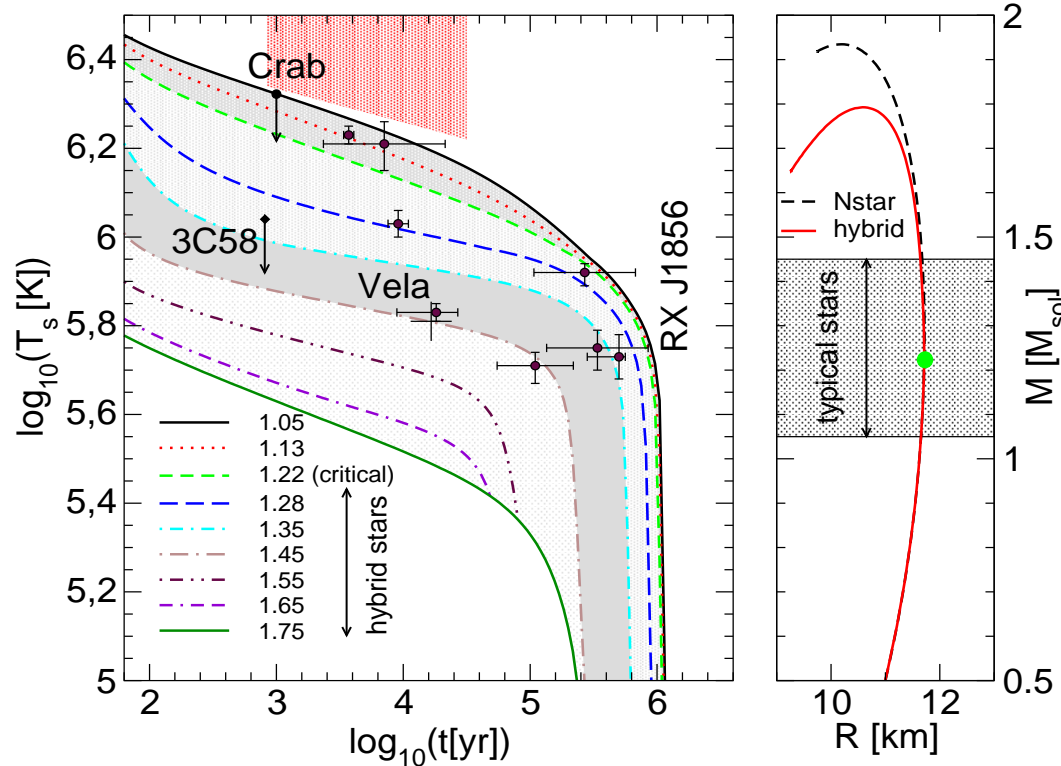


- Low-mass X-ray binary 4U 1636:
 Mass constraint from ISCO obs.



QUARK MATTER IN COMPACT STARS: COOLING CONSTRAINT

Quark matter in compact stars: color superconducting

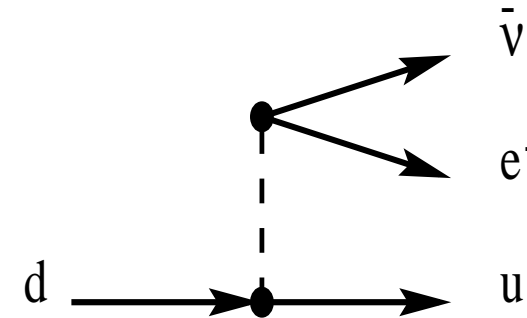


Popov et al: Neutron star cooling constraints ...
 PRC 74, 025803 (2006); [nucl-th/0512098]

- Neutrinos carry energy off the star, Cooling evolution (schematic) by

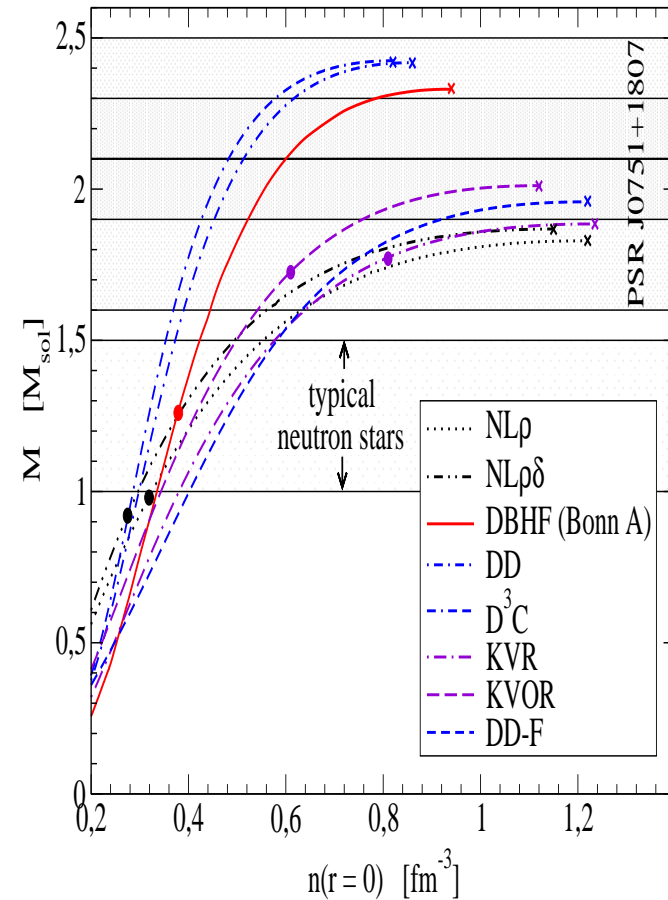
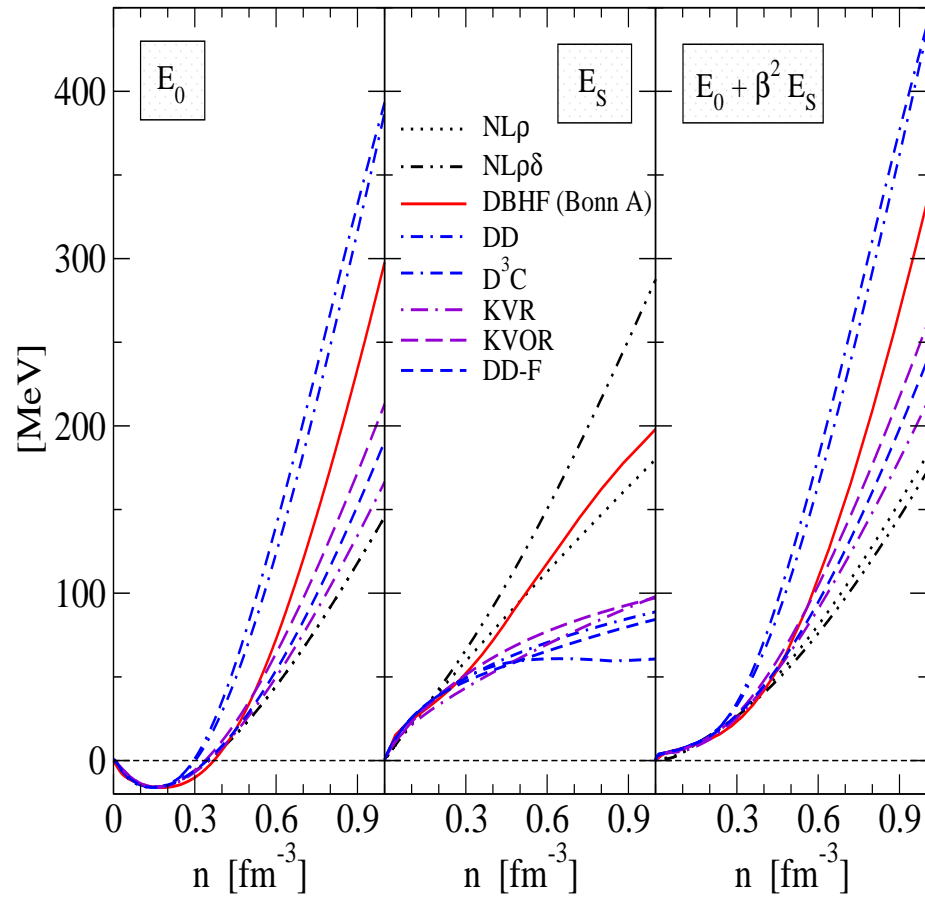
$$\frac{dT(t)}{dt} = - \frac{\epsilon_\gamma + \sum_{j=Urca,\dots} \epsilon_\nu^j}{\sum_{i=q,e,\gamma,\dots} c_V^i}$$

- Most efficient process: Urca



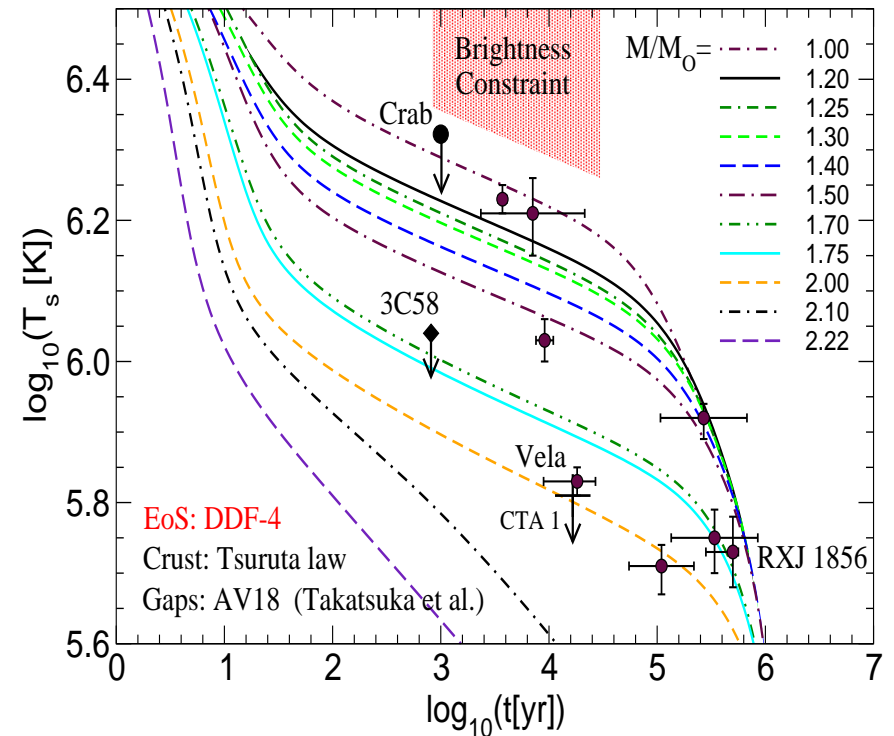
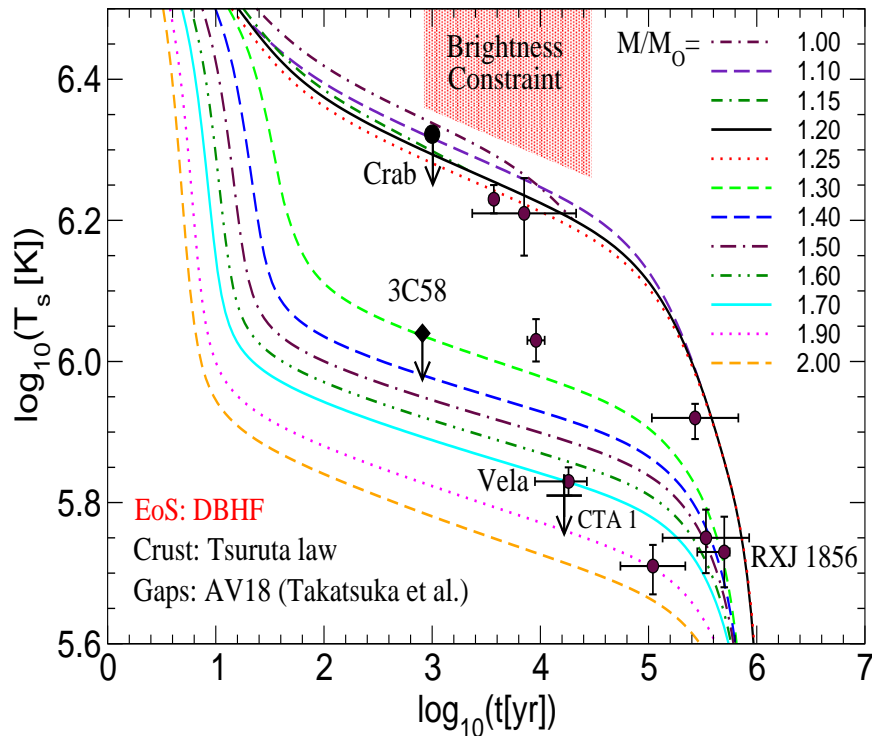
- Exponential suppression by pairing gaps! $\Delta \sim 10 \dots 100 \text{ keV}$

DU THRESHOLD AND 'HADRONIC' NEUTRON STARS



DU threshold for most hadronic EoS active in neutron stars with typical masses !
Klöhn, et al., PRC 74, 035802 (2006); [nucl-th/0602038]

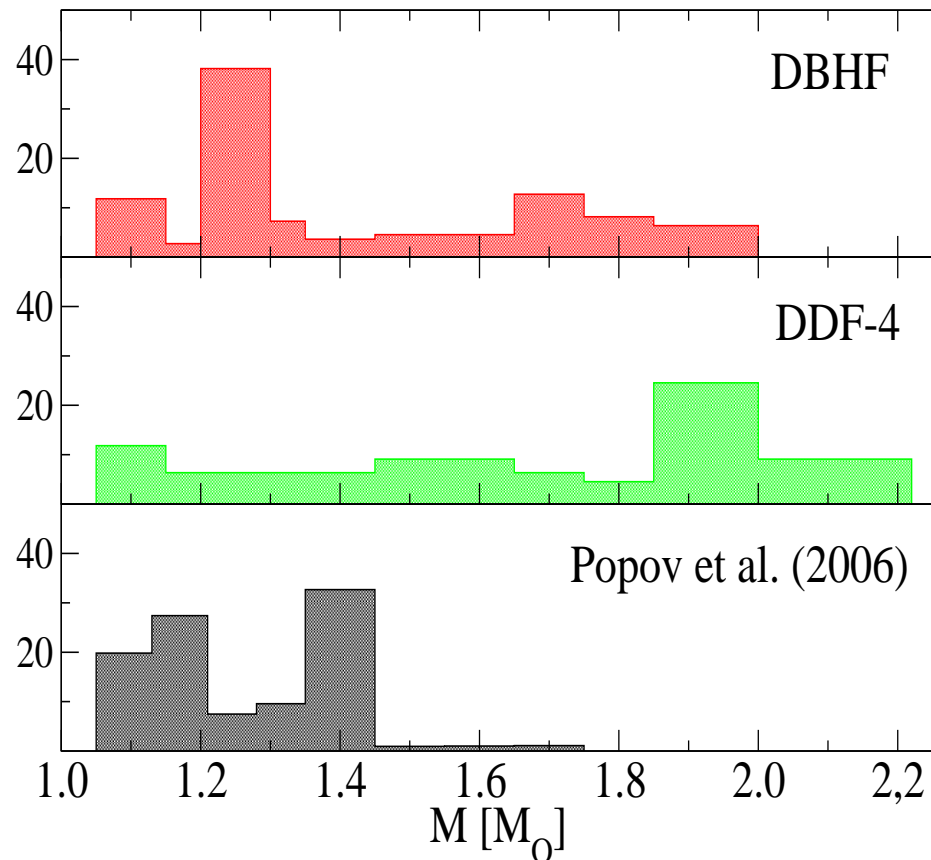
DU THRESHOLD AND 'HADRONIC' NEUTRON STARS (II)



- DU threshold \Rightarrow sensitivity to tiny mass variations;
- Description of Vela not possible with typical masses !

S. Popov et al., PRC 74 (2006); D.B. and H. Grigorian, PPNP (2007) [astro-ph/0612092]

DU THRESHOLD AND 'HADRONIC' NEUTRON STARS (III)

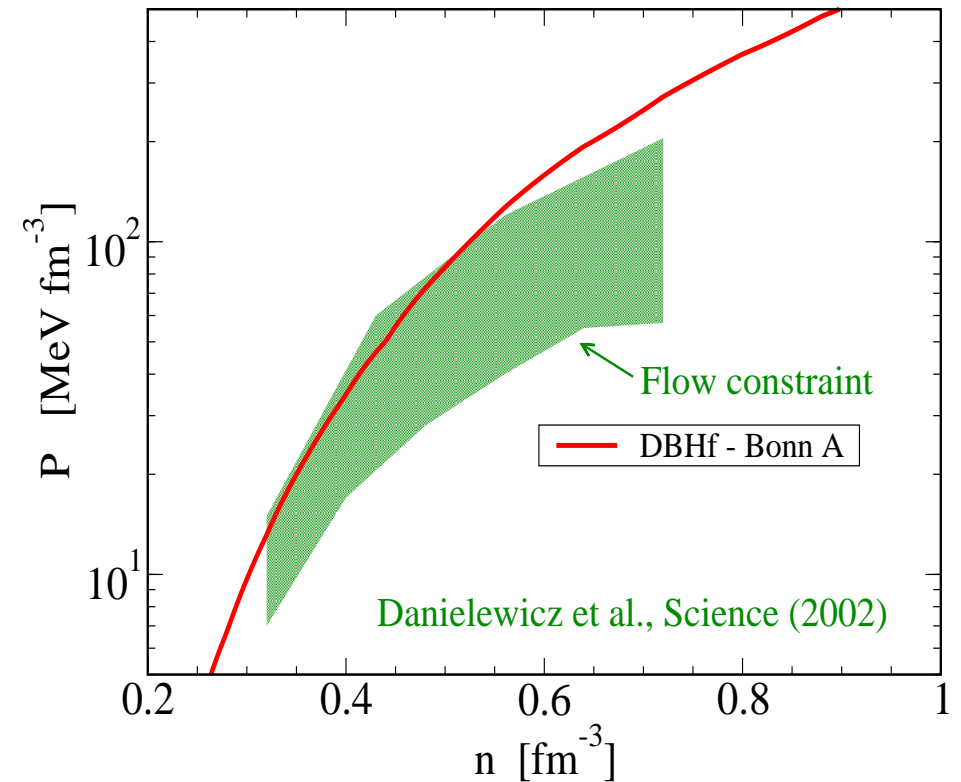
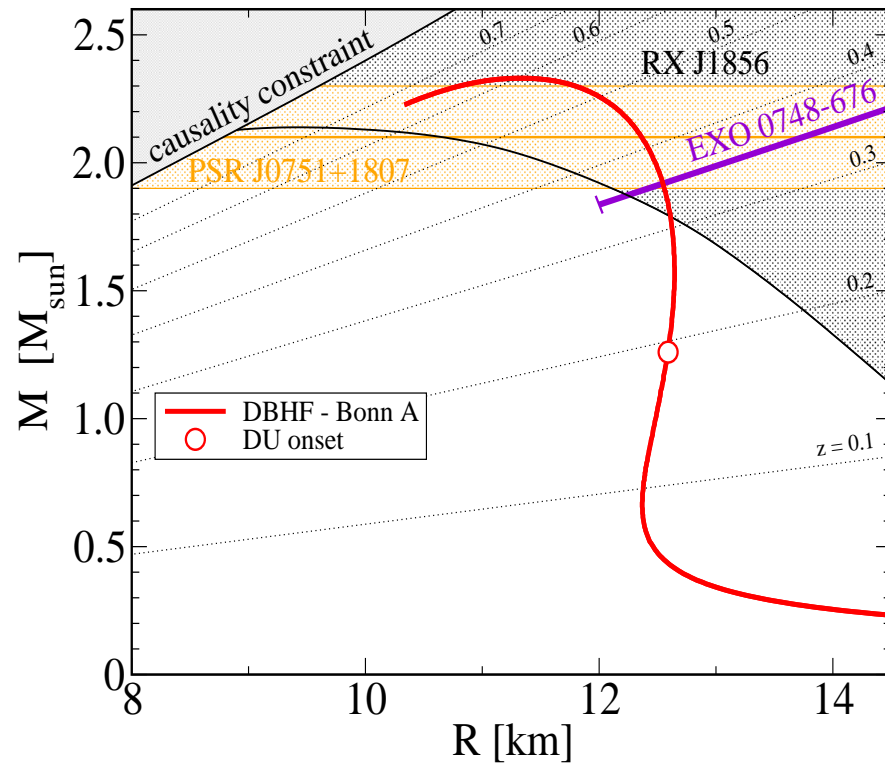


Mass population from Cooling

- DU threshold: overpopulation of a small mass window;
- Hadronic cooling not fast enough to describe Vela with $M < 1.5 M_{\odot}$!

D.B. and H. Grigorian, PPNP (2007);
[astro-ph/0612092]

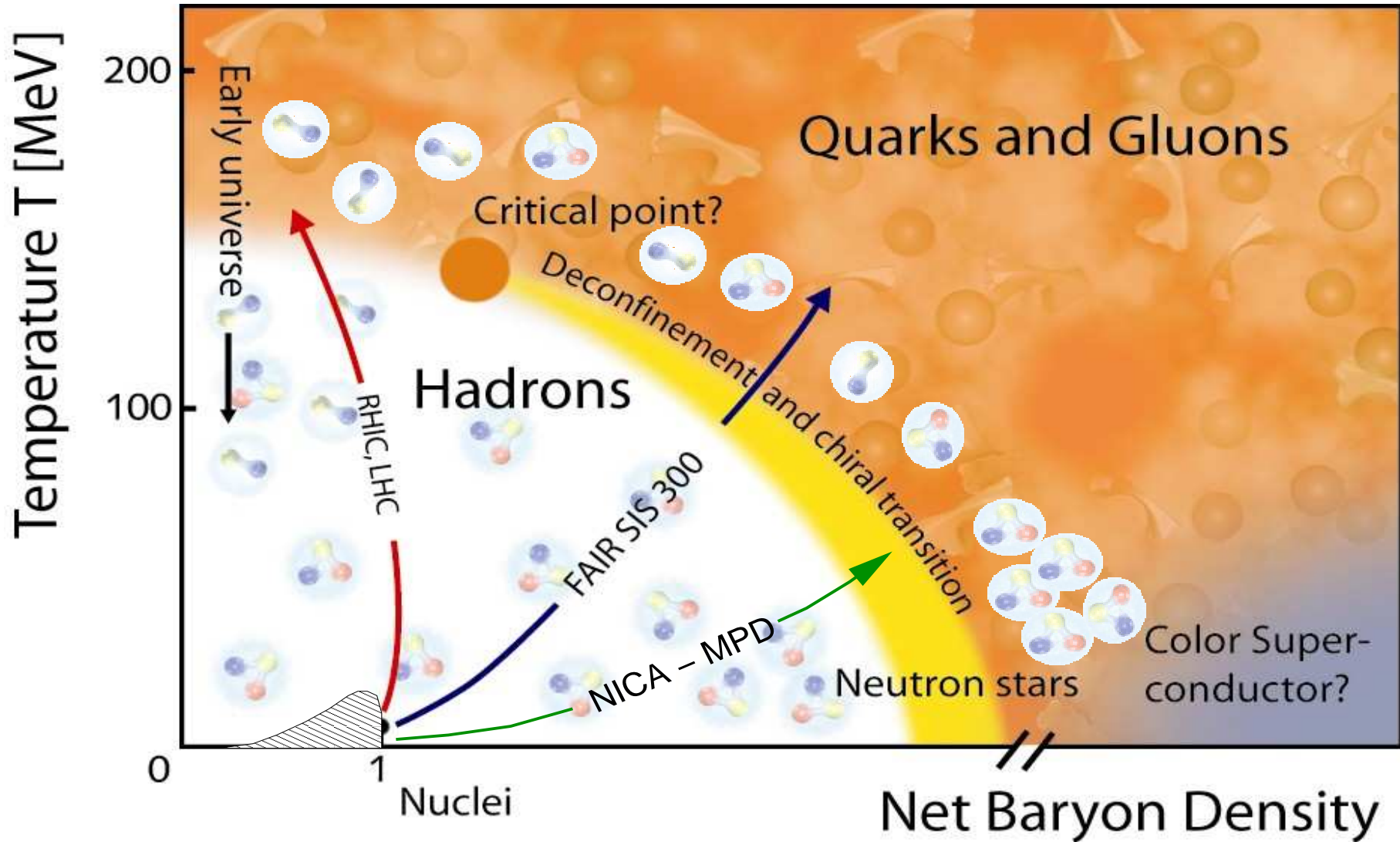
MASS-RADIUS CONSTRAINT AND FLOW CONSTRAINT (I)



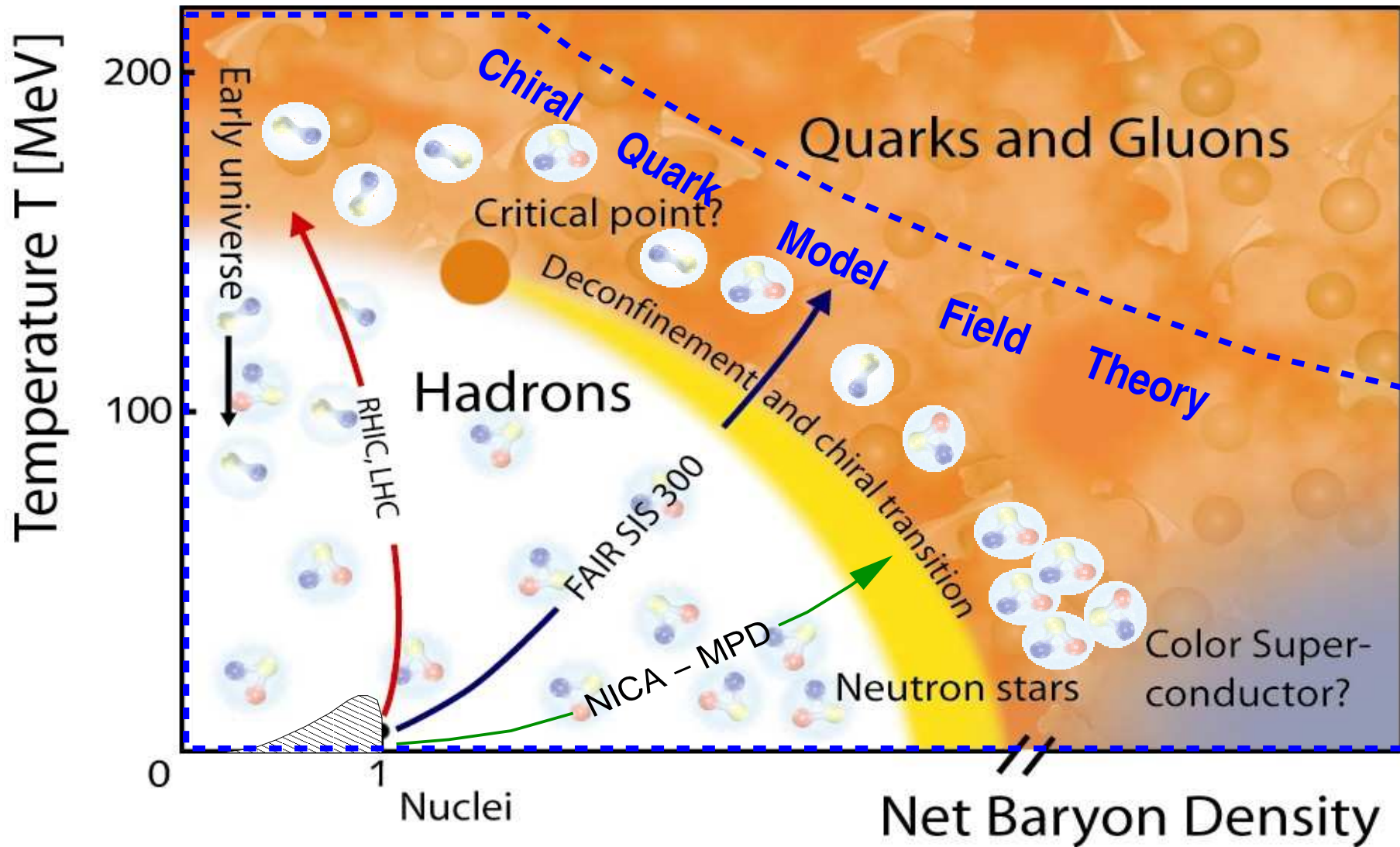
- Large Mass ($\sim 2 M_{\odot}$) and radius ($R \geq 12$ km) \Rightarrow stiff EoS;
But: DBHF has DU onset in typical mass region!
- Flow in Heavy-Ion Collisions \Rightarrow not too stiff EoS ! **But:** DBHF violates at high densities

Klähn, D.B., Sandin, Fuchs, Faessler, Grigorian, Röpke, Trümper, [arxiv:nucl-th/0609067]

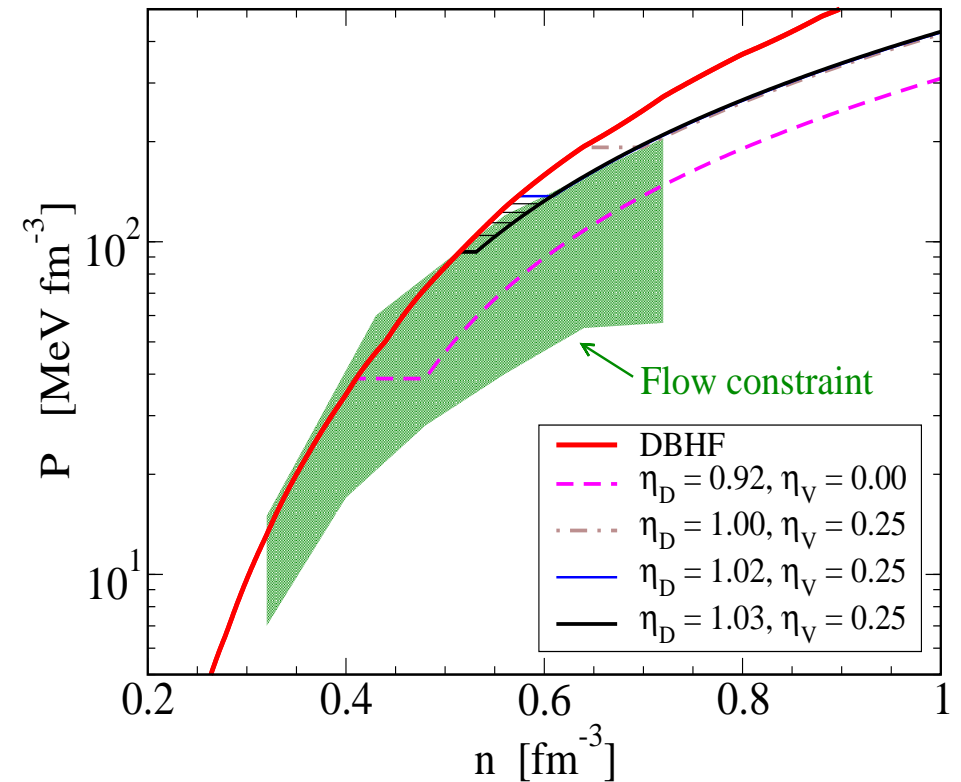
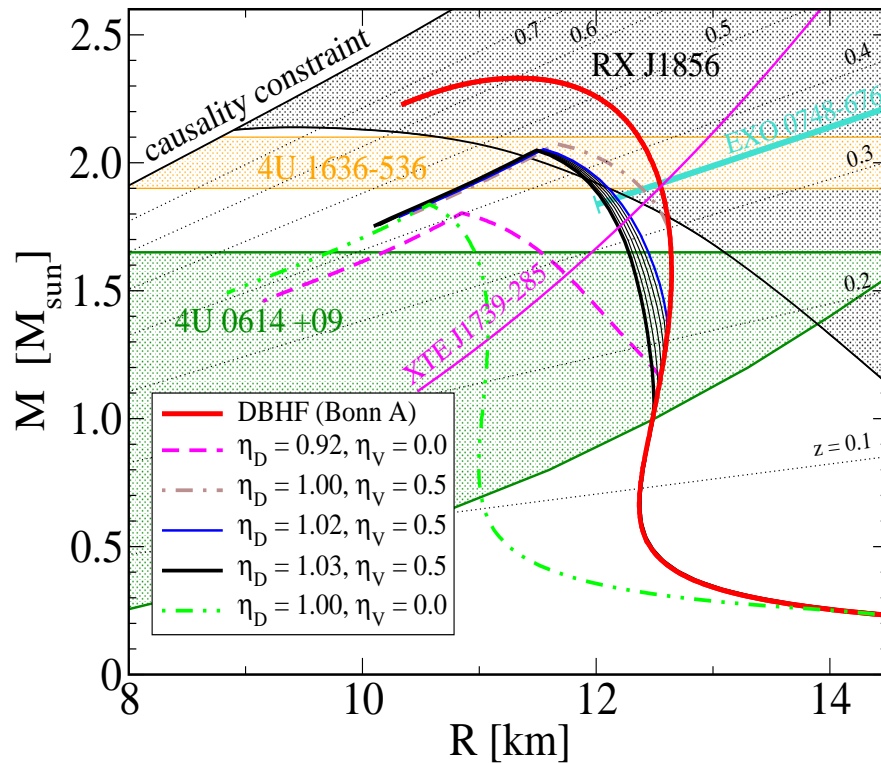
QUARK SUBSTRUCTURE AND PHASE DIAGRAM



PHASE DIAGRAM OF QCD: CHIRAL QUARK MODELS



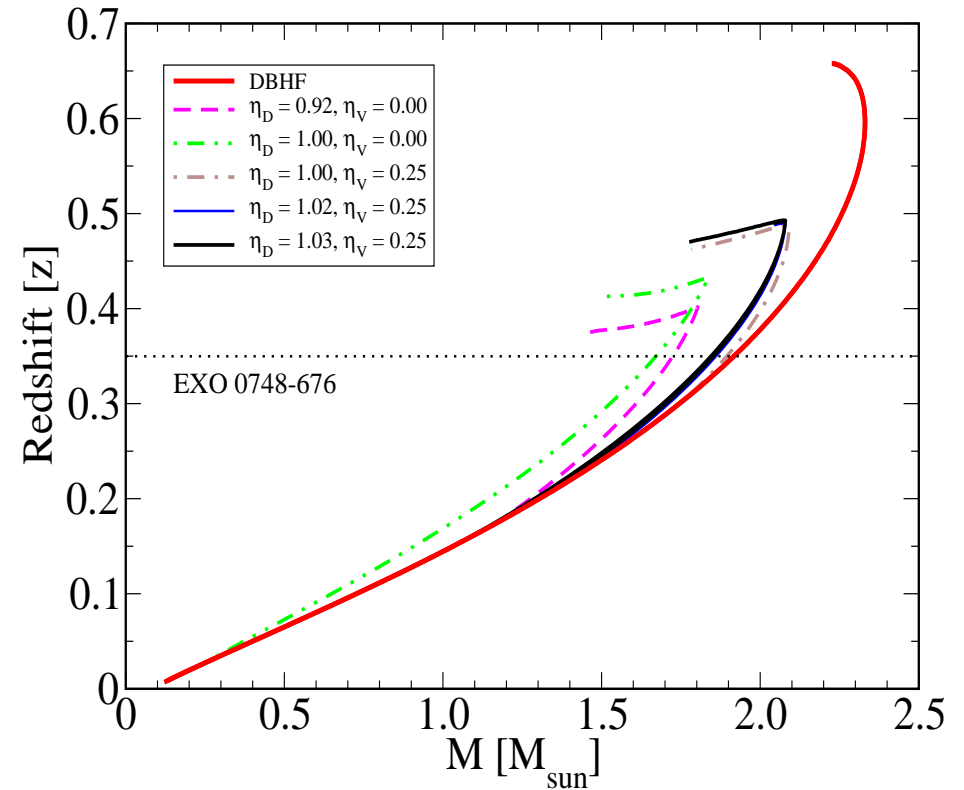
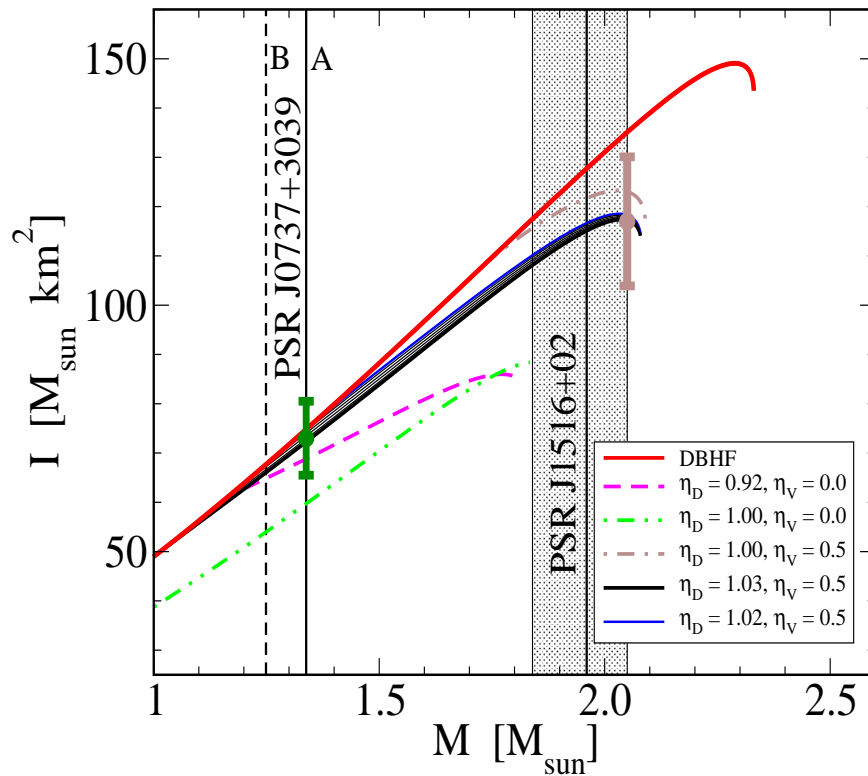
MASS-RADIUS CONSTRAINT AND FLOW CONSTRAINT (II)



- Large Mass ($\sim 2 M_{\odot}$) and radius ($R \geq 12$ km) \Rightarrow stiff quark matter EoS;
 - Note:** DU problem of DBHF removed by deconfinement! **and:** CFL core Hybrids unstable!
- Flow in Heavy-Ion Collisions \Rightarrow not too stiff EoS !
 - Note:** Quark matter removes violation by DBHF at high densities

Klähn, D.B., Sandin, Fuchs, Faessler, Grigorian, Röpke, Trümper, [arxiv:nucl-th/0609067]

HYBRID STARS THAT MASQUERADE AS NEUTRON STARS*



- Moment of Inertia \Rightarrow objects with large masses necessary
- Surface redshift \Rightarrow large values (> 0.5) troublesome for quark matter

* Alford et al., ApJ 629, 969 (2005); Klähn et al., PLB (2007); nucl-th/0609067

HYBRID STAR COOLING WITH 2SC QUARK MATTER

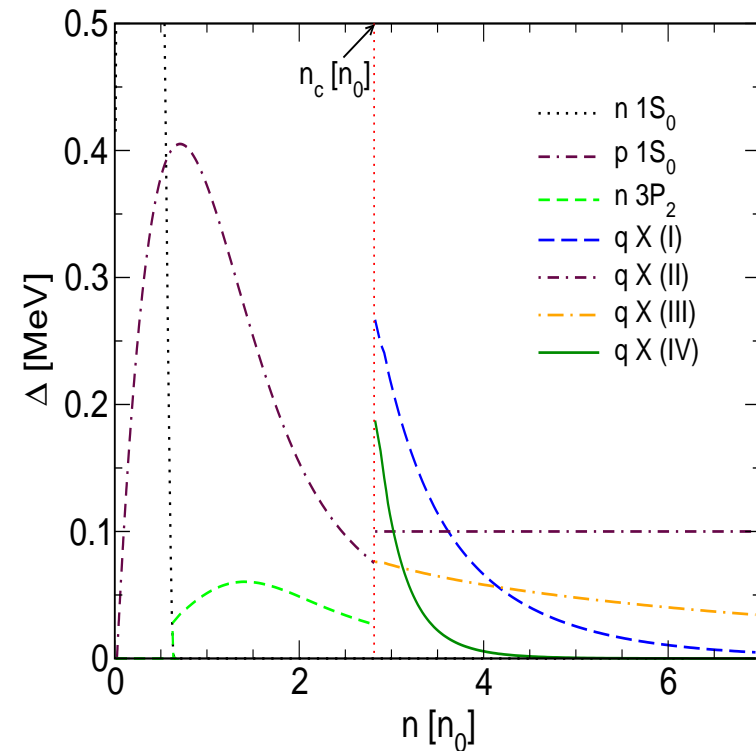
2SC phase: 1 color (blue) is unpaired
(mixed superconductivity)

Ansatz 2SC + X phase:

$$\Delta_X(\mu) = \Delta_0 \exp[\alpha(1 - \mu/\mu_c)]$$

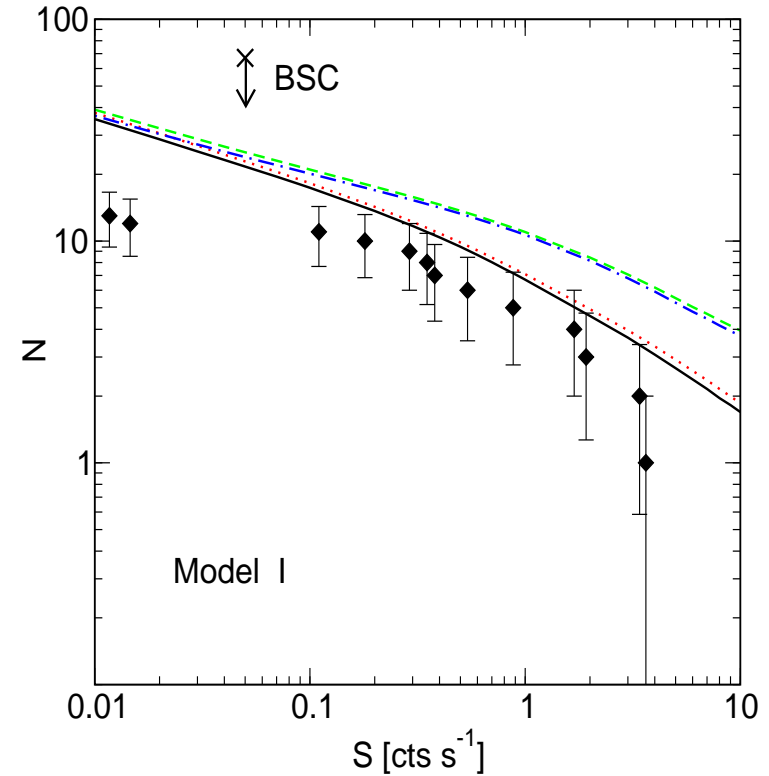
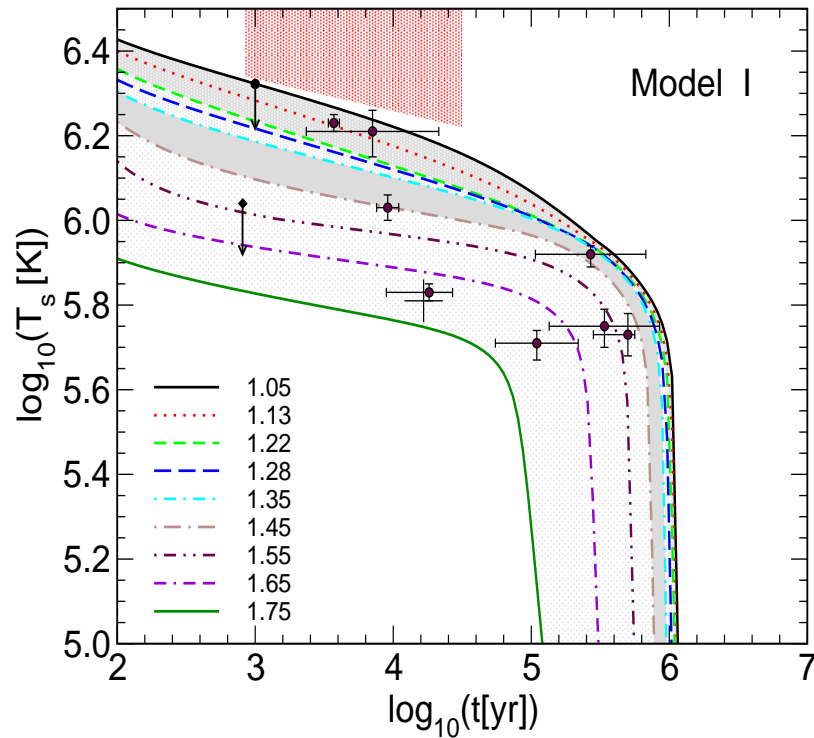
Model	Δ_0 [MeV]	α
I	1	10
II	0.1	0
III	0.1	2
IV	5	25

Popov, Grigorian, D.B., PRC 74 (2006)



Pairing gaps for hadronic phase
(AV18 - Takatsuka et al. (2004))
and 2SC + X phase

HYBRID STAR COOLING WITH 2SC QUARK MATTER (II)

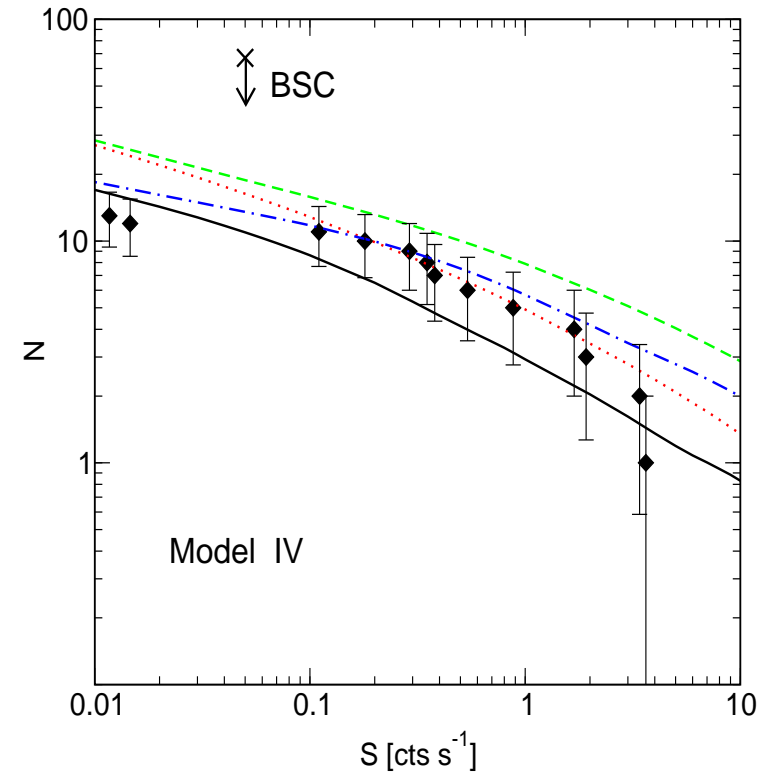
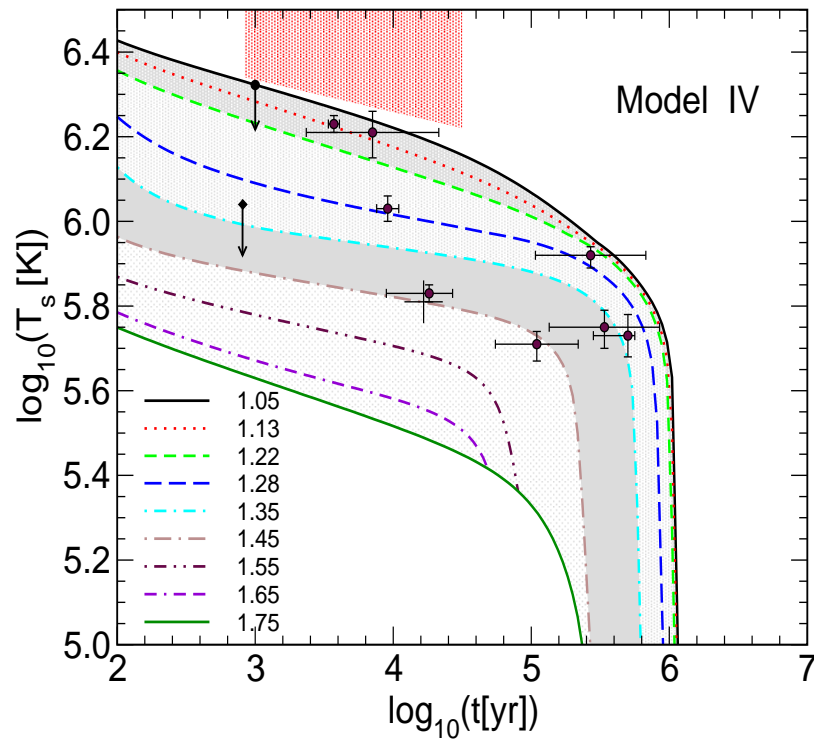


2SC + X phase, $\Delta_0 = 1 \text{ MeV}$, $\alpha = 10$
 Too large mass for Vela required

Popov, Grigorian, D.B., PRC 74 (2006)

Log N - Log S test fails

HYBRID STAR COOLING WITH 2SC QUARK MATTER (III)



2SC + X phase, $\Delta_0 = 5 \text{ MeV}$, $\alpha = 25$
 Temperature-age and Vela mass OK

Log N - Log S test passed

Popov, Grigorian, D.B., PRC 74 (2006)

HYBRID STAR COOLING WITH 2SC QUARK MATTER (IV)

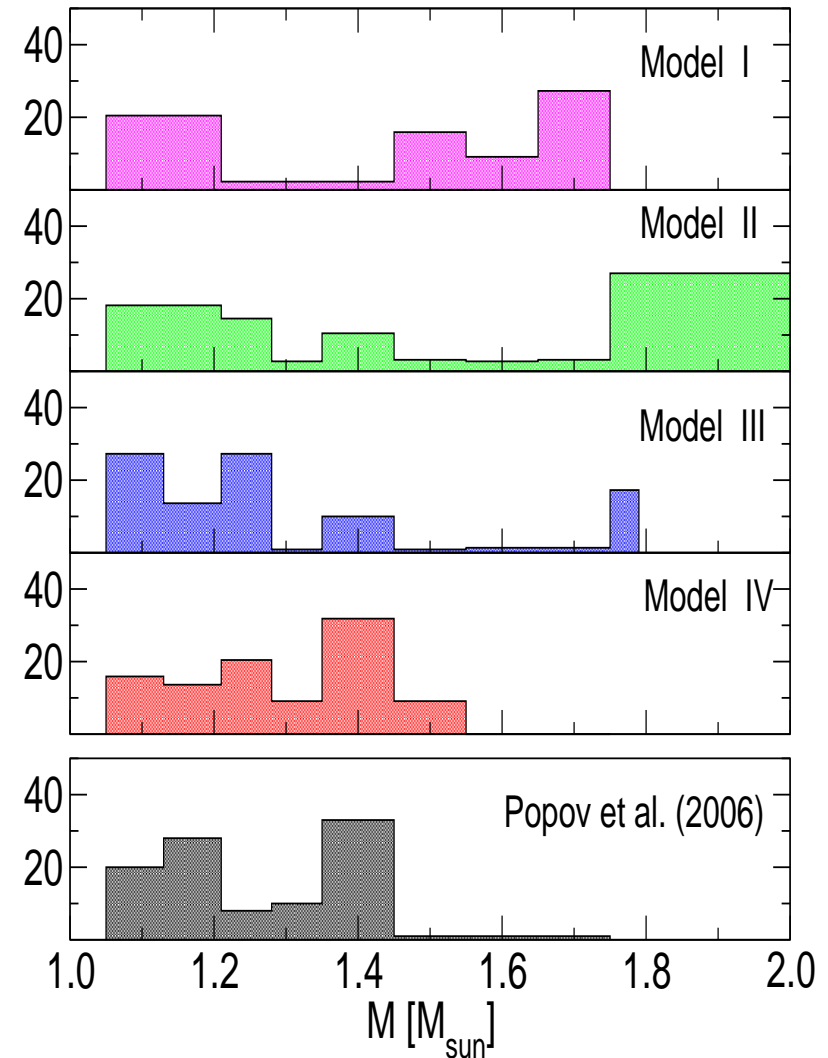
Hybrid star cooling passes all modern tests:

- Temperature - age
- Log N - Log S
- Brightness constraint
- Vela mass (Population synthesis)

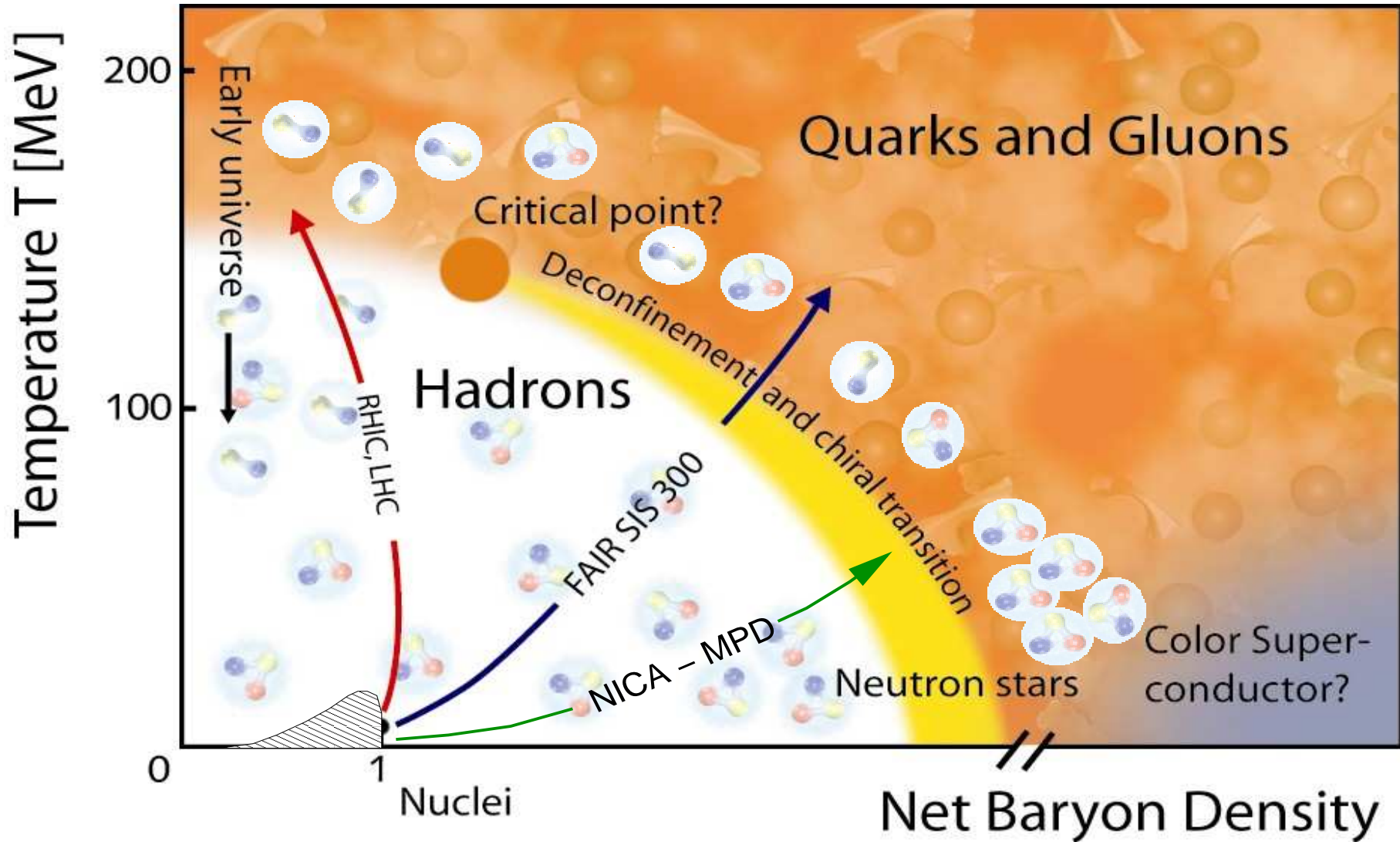
Popov, Grigorian, D.B., PRC 74 (2006)

D.B., H. Grigorian, D.B., PPNP 59 (2007)

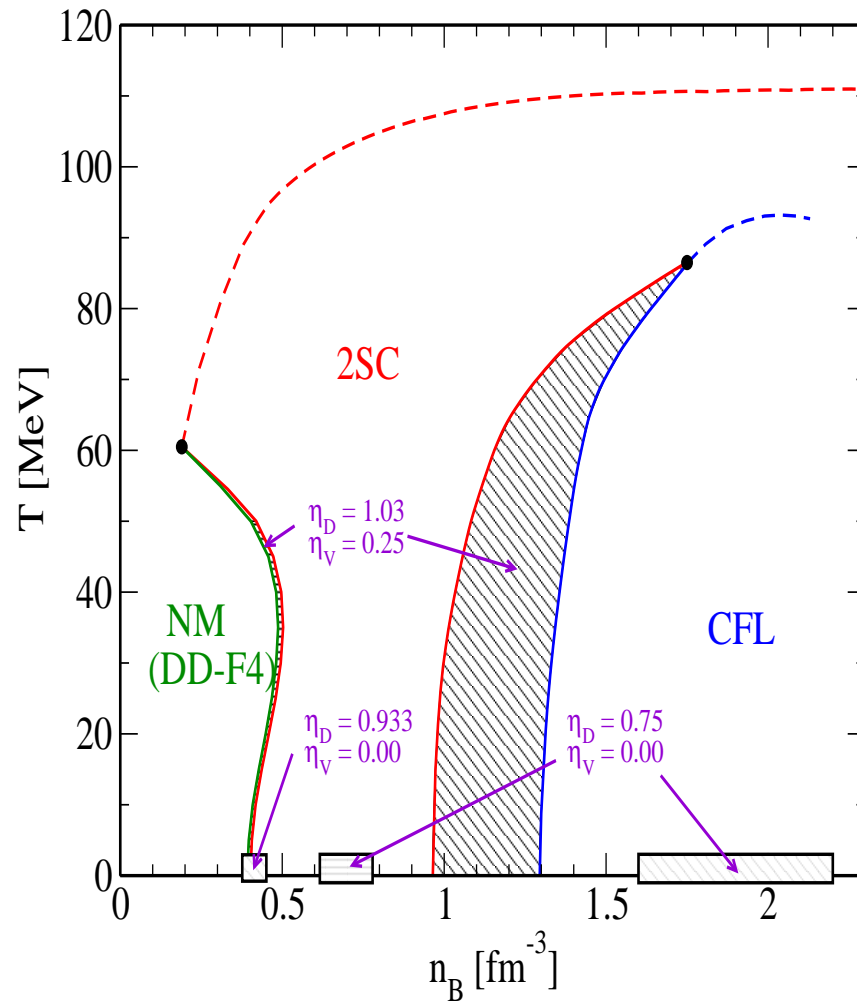
astro-ph/0612092



QUARK SUBSTRUCTURE AND PHASE DIAGRAM



PHASE DIAGRAM FOR SYMMETRIC MATTER (HIC)



- Critical density for deconfinement in heavy-ion collisions $n_{dec} \geq 3 n_0$
- Almost crossover (masquerade!), i.e. no density jump, no latent heat/ time delay!
- 2SC - CFL phase transition at $n \geq 6 n_0$ with density jump and latent heat/ time delay!
Provided the temperature can be kept low $T \leq 70$ MeV

Klahn, Sandin, Typel, D.B., in preparation

SUMMARY

- Mott-Hagedorn model as alternative interpretation of Lattice data
- Microscopic formulation of the hadronic Mott effect within a chiral quark model
- Mesonic (hadronic) correlations important for $T > T_c$ (and $\mu > \mu_c$)
- Step-like enhancement of threshold processes due to Mott effect
- Large neutron star masses: EoS at high densities needs to be stiff! QM - superconducting!
- Hybrid stars 'masquerade' as neutron stars: Cooling observables!

PLANS FOR THE FUTURE

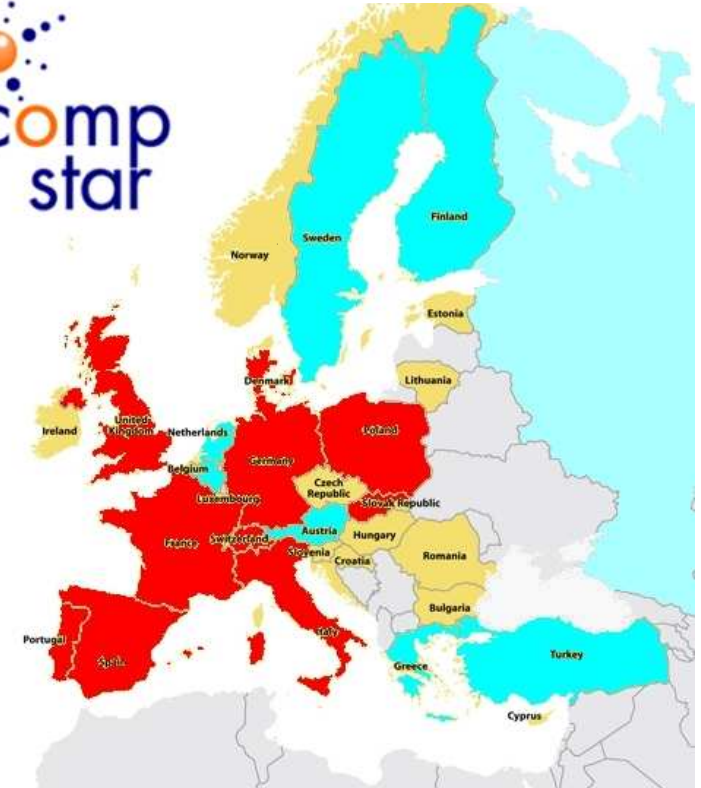
- Bridge Lattice QCD and Phenomenology: spectral functions
- Calculate J/ψ breakup with baryon impact \implies CBM @ FAIR GSI
- Mott effect in dense baryonic matter: nucleon dissociation!
- EoS for hot, dense matter with Mott-effect, encoded in hadronic spectral functions

COLLABORATIONS

EUROPEAN
SCIENCE
FOUNDATION
SETTING SCIENCE AGENDAS FOR EUROPE

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ESF Research Networking Programme
“CompStar”, 2008 - 2013

INVITATIONS ...

DIAS-TH: Dubna International Advanced School of Theoretical Physics
Helmholtz International Summer School

Dense Matter in Heavy Ion Collisions and Astrophysics

Bogoliubov Laboratory of Theoretical Physics
JINR, Dubna, Russia, July 14-26, 2008

TOPICS:

- Hadrons in the Medium
- Equation of state and Phase Transitions
- Hadron Production and Heavy Ion Collisions
- Dense Matter in Compact Stars
- Future Experimental Facilities

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- V. Voronov (JINR)
- D. Blaschke (JINR, U Wroclaw)

LOCAL ORGANIZERS:

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- J. Schmelzer (U Rostock, JINR)
- V. Zhuravlev (JINR)
- V. Skokov (sc. secretary, JINR)
- A. Dolya (secretary, JINR)

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Helmholtz International Summer School

“Dense Matter in Heavy-Ion Collisions and Astrophysics”,

Dubna, Russia, July 14-26, 2008

<http://theor.jinr.ru/~dm2008>

ECT* Workshop

“Quarkonia in Hot and Dense Matter”

Trento, Italy, May 25-29, 2009

<http://www.ect.it>

ESF Research Networking Programme

“CompStar” (2008-2013)

<http://www.esf.org/compstar>

<http://compstar-esf.org>

INVITATION ... TO WROCLAW

